# Siegel-Veech constants in $\mathcal{H}(2)$ 

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#### Abstract

Abelian differentials on Riemann surfaces can be seen as translation surfaces, which are flat surfaces with cone-type singularities. Closed geodesics for the associated flat metrics form cylinders, whose number under a given maximal length was proved by Eskin and Masur to generically have quadratic asymptotics in this length, with a common coefficient for the quadratic asymptotics shared by almost all surfaces in each moduli space of translation surfaces, and called a Siegel-Veech constant.

Square-tiled surfaces are some specific translation surfaces which have their own quadratic asymptotics for the number of cylinders of closed geodesics. It is an interesting question whether the Siegel-Veech constant of a given moduli space can be recovered as a limit of individual constants of square-tiled surfaces in this moduli space. Here we prove that it is the case in the moduli space $\mathcal{H}(2)$ of translation surfaces of genus two with one singularity.


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## 1 Introduction

### 1.1 Geodesics on the torus

On the standard torus $\mathbf{T}^{2}=\mathbf{R}^{2} / \mathbf{Z}^{2}$, the number $N(L)$ of maximal families of parallel simple closed geodesics of length not exceeding $L$ is well-known (and easily seen) to grow quadratically in $L$, with

$$
N(L) \sim \frac{1}{2 \zeta(2)} \cdot \pi L^{2}
$$

which is one half of the asymptotic for the number of primitive lattice points in a disc of radius $L$. The factor one half comes from counting unoriented rather than oriented geodesics.
By convention, the corresponding Siegel-Veech constant is $\frac{1}{2 \zeta(2)}$. Note that it is the coefficient of $\pi L^{2}$, and not of $L^{2}$, in the asymptotic.
Marking the origin of the torus (i.e. artificially considering it as a singularity or saddle), the number of geodesic segments joining the saddle to itself, of length at most $L$, coincides with the number of families of simple closed geodesics.

### 1.2 Geodesics on translation surfaces

It is a standard fact that Abelian differentials on Riemann surfaces can be seen as translation surfaces. On translation surfaces of genus $\geqslant 2$, countings of closed geodesics or saddle connections, similar to those just described for the torus, can be made.

There, the countings of saddle connections and of cylinders of simple closed geodesics do not coincide, but their growth rates remain quadratic. This is made more precise by several related results.

Masur proved [Ma88, Ma90] that for every translation surface, there exist positive constants $c$ and $C$ such that the counting functions of saddle connections and of maximal cylinders of closed geodesics satisfy

$$
c \cdot \pi L^{2} \leqslant N_{\mathrm{cyl}}(L) \leqslant N_{\mathrm{sc}}(L) \leqslant C \cdot \pi L^{2}
$$

for large enough $L$.
Veech [Ve] proved that on a square-tiled surface (and on any Veech surface) there are in fact exact quadratic asymptotics, and Gutkin and Judge [GuJu] gave another proof of that.

Another proof for the upper quadratic bounds for $N_{\text {cyl }}(L)$ and $N_{\mathrm{sc}}(L)$ was given by Vorobets [Vo].

Eskin and Masur [EM] gave yet another one, and proved that for each connected component of each stratum of each moduli space of normalised (i.e., area 1) abelian or quadratic differentials, there are constants $c_{\mathrm{sc}}$ and $c_{\mathrm{cyl}}$ such that almost every surface in the component has $N_{\mathrm{sc}}(L) \sim c_{\mathrm{sc}} \pi L^{2}$ and $N_{\mathrm{cyl}}(L) \sim$ $c_{\mathrm{cyl}} \pi L^{2}$.

It is an interesting open problem whether all translation surfaces have exact quadratic asymptotics for countings of saddle connections and of cylinders of closed geodesics.

The particular constants for many Veech surfaces have been computed explicitly by Veech [Ve], Vorobets [Vo], Gutkin-Judge [GuJu], Schmoll [Schmo]. Constants for some families of non-Veech surfaces were also given by Eskin-MasurSchmoll [EMS] and Eskin-Marklof-Morris [EMWM]. The generic constants for the connected components of all strata of abelian differentials were computed by Eskin, Masur and Zorich in [EMZ].

In general, the particular constants for Veech surfaces do not coincide with the generic constants of the strata where they live.

There is also another subtle difference between Veech surfaces and generic surfaces. Define cylinders as regular if their boundary components both consist of a single saddle connection. In any connected component of stratum in genus $\geqslant 2$, a generic surface has no irregular cylinders, while on Veech surfaces countings of irregular cylinders have quadratic asymptotics.

What we will prove however is that individual 'quadratic constants' for regular cylinders on square-tiled surfaces of the stratum $\mathcal{H}(2)$ (translation surfaces of genus 2 with one singularity) allow to retrieve the generic Siegel-Veech constant of $\mathcal{H}(2)$ as a limit. See Theorem 1 in $\S 1.3$ for a precise statement.

### 1.3 Setting and main result

In this paper, we are concerned with the stratum $\mathcal{H}(2)$ consisting of genus 2 abelian differentials with a double zero, or in other words translation surfaces of genus 2 with one singularity (of angle $6 \pi$ ). We prove:

Theorem 1 Consider a sequence $S_{n}$ of area 1 surfaces in $\mathcal{H}(2)$, each tiled by some prime number $p_{n}$ of squares, with $p_{n} \rightarrow \infty$. Then the constants in the quadratic asymptotics for regular cylinders of closed geodesics on the surfaces $S_{n}$ tend to $\frac{10}{3} \cdot \frac{1}{2 \zeta(2)}$, the Siegel-Veech constant of $\mathcal{H}(2)$ for cylinders of closed geodesics.

Remark It is possible to adapt our calculations to show that the constants in the quadratic asymptotics for irregular cylinders of closed geodesics on the surfaces $S_{n}$ in the theorem tend to 0 , so that the constants in the quadratic asymptotics for all cylinders (both regular and irregular), tend to the generic constant of the stratum $\mathcal{H}(2)$ as well.

Remark We believe that the assumption that the number of squares tiling the surfaces is prime is unnecessary, but we have not yet been able to adapt the calculations to show the convergence of Siegel-Veech constants in the case of nonprime numbers of tiles.

The proof of the theorem relies on fine estimates presented in $\S 3.1$.
Pierre Arnoux pointed out to us the analogy to a result of C. Faivre on Lévy constants of quadratic numbers, see $[F]$ or $[D P]$.

### 1.4 Acknowledgments

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This research was carried out in Montpellier for the most part; some intuition was gained from computer calculations (programmed in Caml Light), run using the Medicis server at École polytechnique.

## 2 Preliminaries

The stratum $\mathcal{H}(2)$ is the simplest stratum of abelian differentials after the (well-understood) stratum of abelian differentials on tori. As every stratum, it admits a natural $\mathrm{SL}(2, \mathbf{R})$ action, and we will recall here some facts concerning the orbits of certain special points of $\mathcal{H}(2)$ : square-tiled surfaces.

A square-tiled surface is a ramified translation cover of the standard torus, with only one branch point. The number of square tiles is the number of sheets of the covering, or the degree of the corresponding covering map to the standard torus. A square-tiled surface is called primitive if this covering map does not factor through a covering of a larger torus with only one branch point.

### 2.1 Orbits of square-tiled surfaces

By a theorem of McMullen [Mc2], in $\mathcal{H}(2)$, for $n>3$, primitive $n$-square-tiled surfaces are all in one $\operatorname{SL}(2, \mathbf{R})$-orbit if $n$ is even, and in exactly two $\operatorname{SL}(2, \mathbf{R})$ orbits if $n$ is odd (see [HL1] for the prime $n$ case). We will denote these orbits by $\mathcal{A}_{n}$ and $\mathcal{B}_{n}$ for odd $n$ and by $\mathcal{E}_{n}$ for even $n$.

The integer points in these orbits are primitive $n$-square-tiled surfaces, and they form $\operatorname{SL}(2, \mathbf{Z})$-orbits which we will denote respectively by $A_{n}, B_{n}, E_{n}$. The number of primitive $n$-square-tiled surfaces in $\mathcal{H}(2)$ is thus the cardinality of $E_{n}$ when $n$ is even and the sum of the cardinalities of $A_{n}$ and $B_{n}$ when $n$ is odd. This number is given in [EMS, Lemma 4.11] to be asymptotic to

$$
\frac{3}{8} n^{3} \prod_{p \mid n}\left(1-\frac{1}{p^{2}}\right)
$$

Formulas for the separate countings of $A_{n}$ and $B_{n}$, conjectured in [HL1], are established in [LeRo] to be

$$
a_{n}=\frac{3}{16}(n-1) n^{2} \prod_{p \mid n}\left(1-\frac{1}{p^{2}}\right), \quad b_{n}=\frac{3}{16}(n-1) n^{2} \prod_{p \mid n}\left(1-\frac{1}{p^{2}}\right)
$$

respectively.
If $n$ tends to infinity within the set of prime numbers, both $a_{n}$ and $b_{n}$ are asymptotic to $\frac{3}{16} n^{3}$.
The natural definition of primitive $n$-square-tiled surfaces gives them area $n$ (each square tile having area 1), but it is sometimes useful to consider the corresponding unit area surfaces, i.e. apply the natural projection from $\mathcal{H}(2)$ to the unit hyperboloid $\mathcal{H}_{1}(2)$.

### 2.2 Cusps

Each square-tiled surface in the stratum $\mathcal{H}(2)$ decomposes into either one or two horizontal cylinders, and can be given as coordinates the heights, widths and twist parameters of these cylinders, see [EMS] or [HL1]. Here we are interested in regular cylinders of closed geodesics, which exist only in twocylinder decompositions (in one-cylinder decompositions, the unique cylinder has three saddle connections on each boundary component).
The decompositions into cylinders provide a way to parametrise square-tiled surfaces (by the heights, widths and twist parameters of their cylinders). These parameters are very convenient to describe the action of $\mathcal{U}=\left\{\left(\begin{array}{cc}1 & n \\ 0 & 1\end{array}\right): n \in \mathbf{Z}\right\}$, which only affects the twist parameters.
The following lemma puts together Lemmas 2.4, 2.5 and 3.1 of [HL1].
Lemma 1 Let $S$ be a primitive $n$-square-tiled surface, and denote by $\mathcal{D}$, resp. $D$, its orbit under $\mathrm{SL}(2, \mathbf{R})$, resp. $\mathrm{SL}(2, \mathbf{Z})$. Then $D$ is the set of primitive $n$ -square-tiled surfaces in $\mathcal{D}$ and the cusps of $\mathcal{D}$ are in bijection with the $\mathcal{U}$-orbits in $D$.

If $S$ has two cylinders, with $h_{i}, w_{i}$ and $t_{i}(i=1,2)$ as height, width and twist parameters, then its cusp width (the cardinality of its $\mathcal{U}$-orbit), is

$$
\operatorname{cw}(S)=\frac{w_{1}}{w_{1} \wedge h_{1}} \vee \frac{w_{2}}{w_{2} \wedge h_{2}} \quad\left(=\frac{w_{1}}{w_{1} \wedge h_{1}} \times \frac{w_{2}}{w_{2} \wedge h_{2}} \text { for prime } n\right)
$$

The surface $S^{\prime}$ with $h_{i}^{\prime}=h_{i}, w_{i}^{\prime}=w_{i}$, and $t_{i}^{\prime}=t_{i} \bmod \left(w_{i} \wedge h_{i}\right)$ is a "canonical" representative of the $\mathcal{U}$-orbit of $S$. Each cusp thus has a unique representative with $0 \leqslant t_{i}^{\prime}<w_{i} \wedge h_{i}$.
(The notations $\wedge$ and $\vee$ are used for gcd and lcm respectively.)
We also recall that given a square-tiled surface $S$, each direction of rational slope on it gives rise to a decomposition of $S$ in cylinders of closed geodesics, and that this direction can be associated to one of the cusps of the $\operatorname{SL}(2, \mathbf{R})$ orbit of $S$.

Note that these cusps can also be understood as cusps of $\Gamma(S)$, the Veech group, or stabiliser under $\operatorname{SL}(2, \mathbf{R})$, of $S$. Algebraically this means conjugacy classes of maximal parabolic subgroups; geometrically the 'cusps' of the quotient surface $\Gamma(S) \backslash \mathbf{H}$.

### 2.3 A formula for the constants

Here, we establish a formula for the constants, for which we will compute estimates in § 3 .

Lemma 2 The number $N_{\text {reg }}(L)$ of regular cylinders of closed geodesics of length $\leqslant L$ on a unit area square-tiled surface $S$ has the following asymptotics:

$$
N_{\mathrm{reg}}(L) \sim \frac{n}{\# D} \sum_{\substack{\mathcal{C}_{j} \text { two-cyl } \\ \text { cusp of } S}} \frac{\mathrm{cw}\left(\mathcal{C}_{j}\right)}{w_{1}^{2}} \frac{1}{2 \zeta(2)} \pi L^{2}
$$

Remark Following a tradition we wrote the asymptotic as a multiple of $\pi L^{2}$ rather than just $L^{2}$, and we wrote $\frac{1}{2 \zeta(2)} \pi L^{2}$ instead of $\frac{3}{\pi} L^{2}$ to bring out the analogy with the corresponding formula for the torus.

Proof We deduce this formula from the material reviewed in $\S \S 2.1-2.2$ above, and from $\S 3$ of $[\mathrm{Ve}]$, to which we refer freely here, both for notations and results.

Veech introduces, for any finite covolume subgroup $\Gamma$ of $\operatorname{PSL}(2, \mathbf{R})$, a complete set $\left\{\Lambda_{j}\right\}_{1 \leqslant j \leqslant r}$ of representatives of the maximal parabolic subgroups of $\Gamma$. We will also refer to the cusps $\mathcal{C}_{j}$. He defines $\Lambda_{0}=\left\{\left(\begin{array}{cc}1 & k \\ 0 & 1\end{array}\right): k \in \mathbf{Z}\right\}$ which we denoted by $\mathcal{U}$.

Then for each $j$ he selects $\beta_{j} \in \operatorname{SL}(2, \mathbf{R})$ which conjugates $\Lambda_{0}$ to $\Lambda_{j}$, i.e. $\beta_{j}^{-1} \Lambda_{j} \beta_{j}=\Lambda_{0}$. When $\Gamma$ is the Veech group of a translation surface $S$, this amounts to represent the cusp $j$ by the surface $\beta_{j}^{-1} S$. Indeed, $\beta_{j} S$ has Veech group $\beta_{j}^{-1} \Gamma \beta_{j}$.

These 'representatives' $\beta_{j}^{-1} S$ of the cusps have width 1 . For square-tiled surfaces, $\Gamma(S)$ is always a subgroup of $\operatorname{SL}(2, \mathbf{Z})$, so it is also usual to conjugate inside $\mathrm{SL}(2, \mathbf{Z})$ to the group generated by some $\left(\begin{array}{cc}1 & k \\ 0 & 1\end{array}\right)$ rather than to $\Lambda_{0}$ itself, thus keeping track of the cusp width (the adequate $k$ ).

Let us illustrate the difference on an example.
Consider the surface $S$, pictured on the left of figure 1, made of seven squares $s_{1}, \ldots, s_{7}$ forming a horizontal cylinder where the right edge of each $s_{i}$ is glued to the left edge of $s_{i+1}$ (indices being understood modulo 7), and where the top edges of squares $s_{1}$ to $s_{7}$ are respectively glued to the bottom edges of squares $s_{3}$ to $s_{6}, s_{1}$ to $s_{2}, s_{7}$.
Consider the direction of the first diagonal. In this direction the surface $S$ decomposes into cylinders of closed geodesics as illustrated on the right of figure 1 (parallel sides of same length identified).


Figure 1:
We get to the standard square-tiled representative of the cusp corresponding to that direction by applying the matrix $M=\left(\begin{array}{cc}1 & 0 \\ -1 & 1\end{array}\right)$ (the matrix in $\operatorname{SL}(2, \mathbf{Z})$ which sends $(1,1)$ to $(1,0)$ and $(0,1)$ to itself). A choice of $\beta_{j}^{-1}$ is $\left(\begin{array}{cc}1 / \sqrt{3} & 0 \\ 0 & \sqrt{3}\end{array}\right) \cdot M$. This sends $3 \cdot(1,1)$ to $(1,0)$.
Figure 2 below represents $\beta_{j}^{-1} S$ on the left, and $M S$ on the right.


Figure 2:
Veech defines $\xi_{j}$ to be the vector $\beta_{j} \cdot(1,0)$. And for each cylinder of closed geodesics in the direction of $\Lambda_{j}$, calling $v$ the holonomy vector for this cylinder, he associates to the cylinder the constant $c_{j}(v)=\frac{\left\|\xi_{j}\right\|}{\|v\|}$.

In our notations, for a surface tiled by unit squares, if $v$ is the holonomy vector of cylinder $i$ of $\operatorname{cusp} \mathcal{C}_{j}$, we have $c_{j}(v)=\sqrt{\operatorname{cw}\left(\mathcal{C}_{j}\right)} / w_{i}$, where $w_{i}$ is the width of this cylinder and $\operatorname{cw}\left(\mathcal{C}_{j}\right)$ the width of this cusp.

Veech's formula for the asymptotics (see [Ve], formula (3.11)) is:

$$
N(L) \sim \operatorname{vol}(\Gamma(S) \backslash \mathbf{H})^{-1}\left(\sum_{j=1}^{r}\left(\sum_{v} c_{j}(v)^{2}\right)\right) L^{2} .
$$

So the contribution of a given cylinder of a cusp $\mathcal{C}_{j}$ to the coefficient of the quadratic asymptotics is $\operatorname{vol}(\Gamma(S) \backslash \mathbf{H})^{-1} c_{j}(v)^{2}$, where $v$ is the holonomy vector of this cylinder.

If we are concerned with regular cylinders of closed geodesics for square-tiled surfaces in $\mathcal{H}(2)$, we need only consider cylinder 1 of two-cylinder cusps.

The volume of the quotient $\Gamma(S) \backslash \mathbf{H}$ equals the index of $\Gamma(S)$ in $\mathrm{SL}(2, \mathbf{Z})$ times the volume of $\mathrm{SL}(2, \mathbf{Z}) \backslash \mathbf{H}$; and the index equals the cardinality of the $\mathrm{SL}(2, \mathbf{Z})$-orbit $D$ of $S$, while the volume of $\mathrm{SL}(2, \mathbf{Z}) \backslash \mathbf{H}$ is $\pi / 3$.

The last thing to observe is the effect of scaling a surface. Consider a surface $S$ where quadratic asymptotics $N(L)=c \cdot \pi L^{2}$ hold, and scale $S$ by a scale factor $r$. On $r S$, the asymptotics become $N(L)=c \cdot \pi(L / r)^{2}$.

A square-tiled surface $S$ of area 1 is a scaled-down version by $1 / \sqrt{n}$ of a surface tiled by $n$ unit squares; this scaling changes the asymptotic by a factor $n$.

This completes the proof of the formula in Lemma 2.

Let us denote by $\widetilde{c}(S)$ the quantity

$$
\frac{n}{\# D} \sum_{\substack{\mathcal{C}_{j} \text { two-cyl } \\ \text { cusp of } S}} \frac{\operatorname{cw}\left(\mathcal{C}_{j}\right)}{w_{1}^{2}}
$$

Our aim is now to prove that $\widetilde{c}(S)$ tends to $\frac{10}{3}$ as the number of square tiles of $S$ tends to infinity staying prime. This will establish Theorem 1.

As a first step for this, using the description of the two-cylinders cusps in the orbits of square-tiled surfaces (see [HL1]), and renaming $w_{1}, w_{2}, h_{1}, h_{2}$ as $a$, $b, h, y$ respectively, we get:

- for $S$ in orbit $A_{n}$ :

$$
\widetilde{c}(S)=\frac{n}{a_{n}}\left(\sum_{\substack{a, b, h, y \geqslant 1 \\ a h+b y=n \\ h \wedge y=1 \\ h, y \operatorname{odd} \\ a<b}} \frac{a b}{a^{2}}+\frac{1}{2} \sum_{\substack{a, b, h, y \geqslant 1 \\ a h+b y=n \\ h \wedge y=1 \\ a \neq b \bmod 2 \\ h \neq y \bmod 2 \\ a<b}} \frac{a b}{a^{2}}\right)
$$

- for $S$ in orbit $B_{n}$ :

$$
\widetilde{c}(S)=\frac{n}{b_{n}}\left(\sum_{\substack{a, b, h, y \geqslant 1 \\ a h+b y=n \\ h \wedge y=1 \\ a, b \text { odd } \\ a<b}} \frac{a b}{a^{2}}+\frac{1}{2} \sum_{\substack{a, b, h, y \geqslant 1 \\ a h+b y=n \\ h \wedge y=1 \\ a \neq b \bmod 2 \\ h \neq y \bmod 2 \\ k<\ell}} \frac{a b}{a^{2}}\right)
$$

The idea is to group two-cylinder cusps sharing the same parameters $w_{1}, w_{2}$, $h_{1}, h_{2}$. Then the sum of the cusp widths adds up to $w_{1} w_{2}$ (for nonprime $n$ some values of the twist parameters could correspond to nonprimitive surfaces, but for prime $n$ all surfaces with $n$ tiles are primitive). All surfaces with $h_{1}$ and $h_{2}$ odd are in orbit $A$, all surfaces with $w_{1}$ and $w_{2}$ odd are in orbit $B$, and those with mixed parities for $w_{i}$ and $h_{i}$ are half in orbit $A$ half in orbit $B$.

## 3 Asymptotics for a large prime number of squares

We need to estimate quantities of the type

$$
\widetilde{c}\left(D_{n}\right)=\frac{n}{\# D_{n}} \sum_{a, b, h, y} \frac{a b}{a^{2}}
$$

where the sum is over positive integers $a, b, h, y$ satisfying certain conditions as above.

### 3.1 A simpler sum

Since $\# D$, for prime $n$, is asymptotically $\frac{3}{16} n^{3}$, we first replace $\frac{n}{\# D_{n}}$ by $\frac{1}{n^{2}}$. Second, we momentarily drop the parity conditions; we will reintroduce them in the following subsections. Last, we drop the condition $a<b$; we will explain later why this does not change the asymptotic.

So we first consider the following simplified sum:

$$
S(n)=\sum_{a \geqslant 1} \frac{1}{a^{2}} \sum_{\substack{b \geqslant 1}} \sum_{\substack{h \geqslant 1, y \geqslant 1 \\ a h+b y=n}} \frac{a b}{n^{2}} .
$$

Denote the sum over $b$ by $S(n, a)$. Introducing the variable $m=b y$,

$$
S(n, a)=\sum_{\substack{1 \leqslant m \leqslant n-a \\ m \equiv n[a]}} \sum_{b \mid m} \frac{a b}{n^{2}}=\frac{a}{n^{2}} \cdot F(n-a, n, a)
$$

where

$$
F(x, k, q)=\sum_{\substack{1 \leqslant m \leqslant x \\ m \equiv k[q]}} \sum_{b \mid m} b .
$$

The following asymptotics hold for $F(x, k, q), S(n, a)$ and $S(n)$.

Lemma 3 For $k \wedge q=1$, and $x \rightarrow \infty$,

$$
F(x, k, q)=\frac{x^{2}}{q} \cdot \frac{\pi^{2}}{12} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O_{q}(x \log x)
$$

## Lemma 4

$$
S(n, a) \xrightarrow[\substack{n \rightarrow \infty \\ n \text { prime }}]{ } \frac{\pi^{2}}{12} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)
$$

Lemma $5 S(n) \xrightarrow[\substack{n \rightarrow \infty \\ n \text { prime }}]{ } \frac{5}{4}$.

Proof of Lemma 3 If $m$ is prime to $q$, denote by $\bar{m}$ the integer in $\{0, \ldots, q-$ $1\}$ such that $\bar{m} m \equiv 1[q]$, and by $u=u(m, k, q)$ the integer in $\{0, \ldots, q-1\}$ such that $u \equiv \bar{m} k[q]$; error terms depend on $q$.

$$
\begin{aligned}
& F(x, k, q)=\sum_{\substack{1 \leqslant m d \leqslant x \\
m d \equiv k[q]}} d \\
& =\sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}} \sum_{\substack{1 \leqslant d \leqslant x / m \\
d \equiv \bar{m} k[q]}} d \\
& =\sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}} \sum_{\substack{1 \leqslant d \leqslant x / m \\
d \equiv u[q]}} d \\
& =\sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}} \sum_{1 \leqslant u+\lambda q \leqslant x / m}(u+\lambda q) \\
& =\sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}}\left(\left(\sum_{1 \leqslant \lambda \leqslant \frac{1}{q}\left(\frac{x}{m}-u\right)} \lambda q\right)+O\left(\frac{x}{m}\right)\right) \\
& =\sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}}\left(\frac{1}{2} q\left(\frac{x}{q m}\right)^{2}+O\left(\frac{x}{m}\right)+O(1)\right) \\
& =\frac{x^{2}}{2 q} \sum_{\substack{1 \leqslant m \leqslant x \\
m \wedge q=1}} \frac{1}{m^{2}}+O(x \log x)
\end{aligned}
$$

To sum only over the integers $m$ with $m \wedge q=1$, we can sum over all $m$ with a factor $\mu(m \wedge q)$, so that all terms cancel out except the ones we want.

$$
\begin{aligned}
F(x, k, q) & =\frac{x^{2}}{2 q} \sum_{d \mid q}\left(\frac{\mu(d)}{d^{2}} \sum_{m \leqslant x / d} \frac{1}{m^{2}}\right)+O(x \log x) \\
& =\frac{x^{2}}{2 q} \sum_{d \mid q} \frac{\mu(d)}{d^{2}}\left(\frac{\pi^{2}}{6}+O(1 / x)\right)+O(x \log x) \\
& =\frac{x^{2}}{q} \cdot \frac{\pi^{2}}{12} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) .
\end{aligned}
$$

Proof of Lemma 4 Lemma 4 follows immediately from Lemma 3 by a dominated convergence argument (similar arguments were used in [HL1, §7]).

Proof of Lemma 5 Lemma 5 is a consequence of Lemma 4 by the following observation.

$$
\begin{aligned}
\sum_{a \geqslant 1} \frac{1}{a^{2}} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right) & =\prod_{p}\left(1+\sum_{\nu \geqslant 1} p^{-2 \nu}\left(1-p^{-2 \nu}\right)\right)=\prod_{p}\left(1+p^{-2}\right) \\
& =\prod_{p} \frac{1-p^{-4}}{1-p^{-2}}=\frac{\zeta(2)}{\zeta(4)}=\frac{\pi^{2} / 6}{\pi^{4} / 90}=\frac{15}{\pi^{2}}
\end{aligned}
$$

### 3.2 Sums with specified parities

We introduce sub-sums of $S(n)$ for specified parities of the parameters.
The observation we just made will need to be completed by the following one.

$$
\sum_{\substack{a \geqslant 1 \\ a \text { even }}} \frac{1}{a^{2}} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)=\sum_{\substack{a \geqslant 1 \\ a \text { odd }}} \frac{1}{4 a^{2}} \frac{3}{4} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)+\sum_{\substack{a \geqslant 1 \\ a \text { even }}} \frac{1}{4 a^{2}} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)
$$

so that $\sum_{\substack{a \geqslant 1 \\ a \text { odd }}} \frac{1}{a^{2}} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)=\frac{12}{\pi^{2}}$ and $\sum_{\substack{a \geqslant 1 \\ a \text { even }}} \frac{1}{a^{2}} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)=\frac{3}{\pi^{2}}$.

### 3.2.1 Odd widths

We now consider the sum over odd $a$ and $b$ :

$$
S^{\mathrm{ow}}(n)=\sum_{\substack{a \geqslant 1 \\ a \text { odd }}} \frac{1}{a^{2}} \sum_{\substack{b \geqslant 1 \\ b \text { odd }\\}} \sum_{\substack{h \geqslant 1, y \geqslant 1 \\ a h+b y=n}} \frac{a b}{n^{2}}
$$

We proceed as for the sum $S(n)$ : putting

$$
\begin{gathered}
F^{\mathrm{ow}}(x, k, q)=\sum_{\substack{1 \leqslant m \leqslant x \\
m \equiv k[q] b \text { odd }}} \sum_{b \mid m} b \text { and } S^{\mathrm{ow}}(n, a)=\frac{a}{n^{2}} \cdot F^{\mathrm{ow}}(n-a, n, a) \\
S^{\mathrm{ow}}(n)=\sum_{\substack{a \geqslant 1 \\
a \text { odd }}} \frac{1}{a^{2}} S^{\mathrm{ow}}(n, a)
\end{gathered}
$$

The following asymptotics hold for $F^{\text {ow }}(x, k, q), S^{\text {ow }}(n, a)$ and $S^{\text {ow }}(n)$.

Lemma 6 For odd $q$, odd $k$, and $x \rightarrow \infty$,

$$
F^{o w}(x, k, q)=\frac{x^{2}}{q} \frac{\pi^{2}}{24} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x)
$$

For odd a,

$$
S^{o w}(n, a) \underset{\substack{n \rightarrow \infty \\ n \text { prime }}}{ } \frac{\pi^{2}}{24} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right)
$$

Finally,

$$
S^{o w}(n) \xrightarrow[\substack{n \rightarrow \infty \\ n \text { prime }}]{ } \frac{1}{2} .
$$

## Proof

$$
\begin{aligned}
F^{\mathrm{ow}}(x, k, q) & =\sum_{t \geqslant 0} \sum_{\substack{1 \leqslant m \leqslant x / 2^{t} \\
2^{t} m \equiv k[q] \\
m \equiv 1[2]}} \sum_{b \mid m} b \\
& =\sum_{t \geqslant 0}\left(\frac{\left(x / 2^{t}\right)^{2}}{2 q} \frac{\pi^{2}}{12} \prod_{p \mid 2 q}\left(1-\frac{1}{p^{2}}\right)+O\left(\left(x / 2^{t}\right) \log \left(x / 2^{t}\right)\right)\right) \\
& =\frac{x^{2}}{q} \frac{1}{1-\frac{1}{4}} \frac{\pi^{2}}{24}\left(1-\frac{1}{2^{2}}\right) \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) \\
& =\frac{x^{2}}{q} \frac{\pi^{2}}{24} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x)
\end{aligned}
$$

### 3.2.2 Odd heights

We now consider the sum over odd $h$ and $y$ :

$$
S^{\mathrm{oh}}(n)=\sum_{a \geqslant 1} \frac{1}{a^{2}} \sum_{\substack{b \geqslant 1}} \sum_{\substack{h \geqslant 1, y \geqslant 1 \\ h, y \text { odd } \\ a h+b y=n}} \frac{a b}{n^{2}} .
$$

Proceeding as previously, we are led to introduce

$$
\begin{aligned}
F^{\mathrm{oh}}(x, k, q)= & \sum_{\substack{1 \leqslant m \leqslant x \\
m \equiv k+q[2 q]]}} \sum_{\substack{b \mid m \\
m \text { odd }}} b \text { and } S^{\mathrm{oh}}(n, a)=\frac{a}{n^{2}} \cdot F^{\mathrm{oh}}(n-a, n, a), \\
& \text { and to write } S^{\mathrm{oh}}(n)=\sum_{a \geqslant 1} \frac{1}{a^{2}} S^{\mathrm{oh}}(n, a) .
\end{aligned}
$$

The following asymptotics hold for $F^{\text {oh }}(x, k, q), S^{\text {oh }}(n, a)$ and $S^{\text {oh }}(n)$.
Lemma 7 For even $q$, odd $k$, and $x \rightarrow \infty$,

$$
F^{o h}(x, k, q)=\frac{x^{2}}{q} \frac{\pi^{2}}{24} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) .
$$

For odd $q$, odd $k$, and $x \rightarrow \infty$,

$$
F^{o h}(x, k, q)=\frac{x^{2}}{q} \frac{\pi^{2}}{32} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) .
$$

For even $a$,

$$
S^{\mathrm{oh}}(n, a) \underset{\substack{n \rightarrow \infty \\ n \text { prime }}}{ } \frac{\pi^{2}}{24} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right) .
$$

For odd $a$,

$$
S^{\mathrm{oh}}(n, a) \underset{\substack{n \rightarrow \infty \\ n \text { prime }}}{ } \frac{\pi^{2}}{32} \prod_{p \mid a}\left(1-\frac{1}{p^{2}}\right) .
$$

Finally,

$$
S^{o h}(n) \underset{\substack{n \rightarrow \infty \\ n \text { prime }}}{ } \frac{1}{2} .
$$

Proof For even $q$ and odd $k$ :

$$
\begin{aligned}
F^{\mathrm{oh}}(x, k, q) & =\sum_{\substack{1 \leqslant m \leqslant x \\
m \equiv k+q}} \sum_{b \mid m]} b \\
& =\frac{x^{2}}{2 q} \frac{\pi^{2}}{12} \prod_{p \mid 2 q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) \\
& =\frac{x^{2}}{q} \frac{\pi^{2}}{24} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) .
\end{aligned}
$$

For odd $q$ and odd $k$ :

$$
\begin{aligned}
F^{\mathrm{oh}}(x, k, q) & =\sum_{t \geqslant 1} \sum_{\substack{1 \leqslant m \leqslant x / 2^{t} \\
2^{t} m \equiv k+q[2 q] \\
m \text { odd }}} \sum_{b \mid m} 2^{t} b \\
& =\sum_{t \geqslant 1} 2^{t} \sum_{\substack{1 \leqslant m \leqslant x / 2^{t} \\
2^{t-1} m \equiv \frac{k+q}{2}[q] \\
m \text { odd }}} \sum_{b \mid m} b \\
& =\sum_{t \geqslant 1} 2^{t} \frac{\left(x / 2^{t}\right)^{2}}{2 q} \frac{\pi^{2}}{12} \prod_{p \mid 2 q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) \\
& =\sum_{t \geqslant 1} \frac{1}{2^{t}} \frac{x^{2}}{q} \frac{\pi^{2}}{24}\left(1-\frac{1}{2^{2}}\right) \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) \\
& =\frac{x^{2}}{q} \frac{\pi^{2}}{32} \prod_{p \mid q}\left(1-\frac{1}{p^{2}}\right)+O(x \log x) .
\end{aligned}
$$

### 3.2.3 Mixed parities

Dealing with the even-odd sums as above would be most cumbersome; this is fortunately not necessary. Indeed, since $S(n)=S^{\mathrm{ow}}(n)+S^{\mathrm{oh}}(n)+S^{\mathrm{eo}}(n)$, and we know the limits of $S(n), S^{\text {ow }}(n)$ and $S^{\text {oh }}(n)$ when $n$ tends to infinity staying prime, we have:

$$
S^{\mathrm{eo}}(n) \xrightarrow[\substack{n \rightarrow \infty \\ n \text { prime }}]{ } \frac{1}{4}
$$

### 3.3 Asymptotics for orbits A and B

We end by showing that the limit we obtained is unchanged by adding a condition $a<b$.
Indeed, since $\#\{(h, y): h \geqslant 1, y \geqslant 1, a h+b y=n\} \leqslant n$, the sum $\sum_{\substack{b=1}}^{a} \sum_{\substack{h \geqslant 1, y \geqslant 1 \\ a h+b y=n}} \frac{a b}{n^{2}}$
is $O(1 / n)$, where the constant of the $O$ depends on $a$.
Putting things together, $\widetilde{c}\left(A_{n}\right)$ and $\widetilde{c}\left(B_{n}\right)$ have the same asymptotics: $S^{\mathrm{A}}(n)=$ $\frac{16}{3}\left(S^{\mathrm{oh}}(n)+\frac{1}{2} S^{\mathrm{eo}}(n)\right)$ and $S^{\mathrm{B}}(n)=\frac{16}{3}\left(S^{\mathrm{ow}}(n)+\frac{1}{2} S^{\mathrm{eo}}(n)\right)$, so they both tend to $\frac{10}{3}$ as $n$ tends to infinity, $n$ prime.

## 4 Concluding remarks

Numerical evidence suggests that the convergence to the generic constant of the stratum occurs not only for prime $n$ but for general $n$; however a proof would involve some complications in the calculations which would make the exposition tedious.

A similar study for the constants that appear in the quadratic asymptotics for the countings of saddle connections could also be made. There one has to take into consideration both one-cylinder and two-cylinder cusps, and some interesting phenomena can be observed: numerical calculations suggest that the sum of the contributions of one-cylinder and two-cylinder cusps has a limit, but separate countings for one-cylinder cusps do not have a limit for general $n$; their asymptotics have fluctuations involving the prime factors of $n$.

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