



## Evaluation of a multiphase LB model for coating flows

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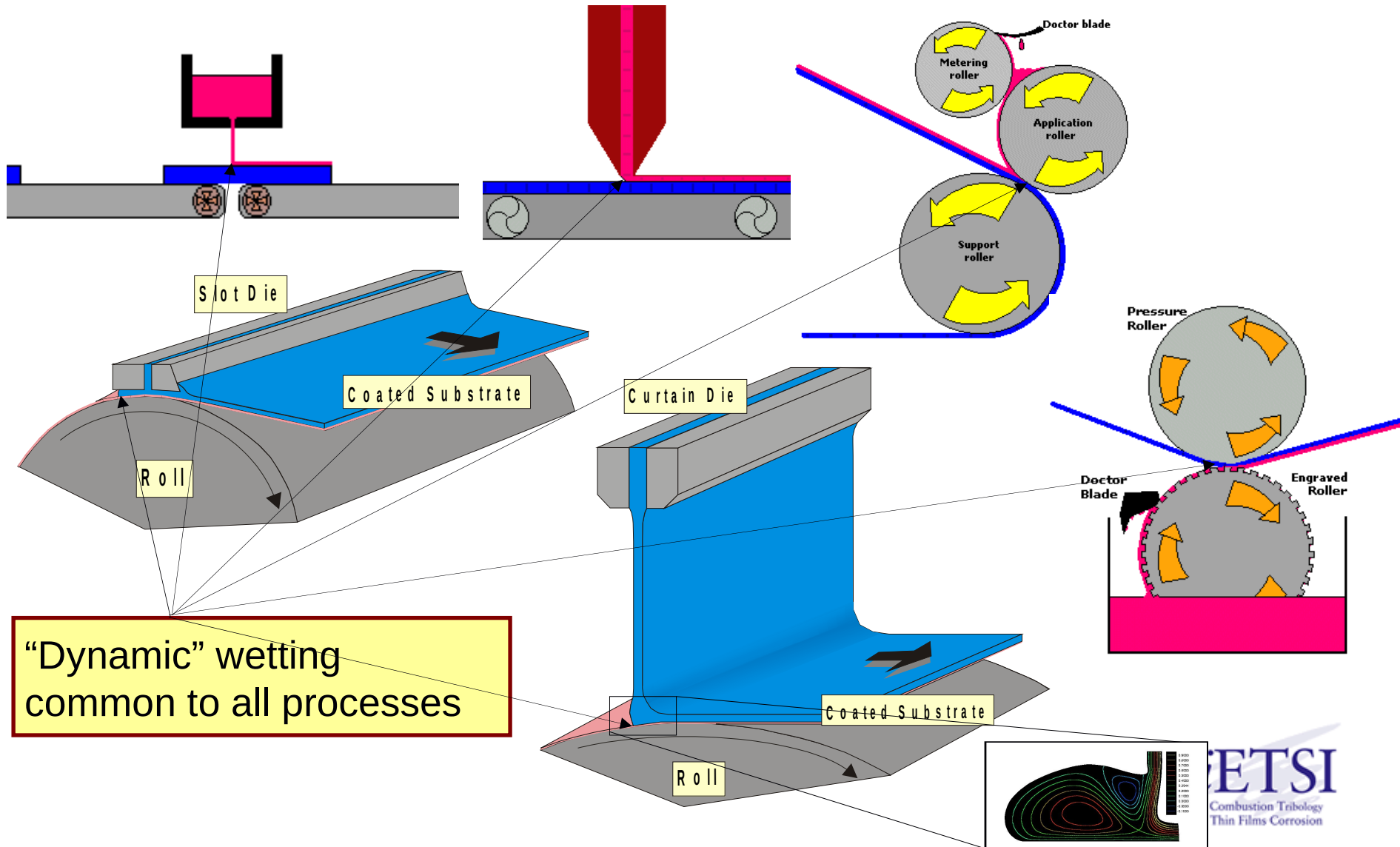
Schéma de Boltzmann sur réseau ; méthodes et applications,  
CEMAGRAF, Paris. Vendredi 05 décembre 2008



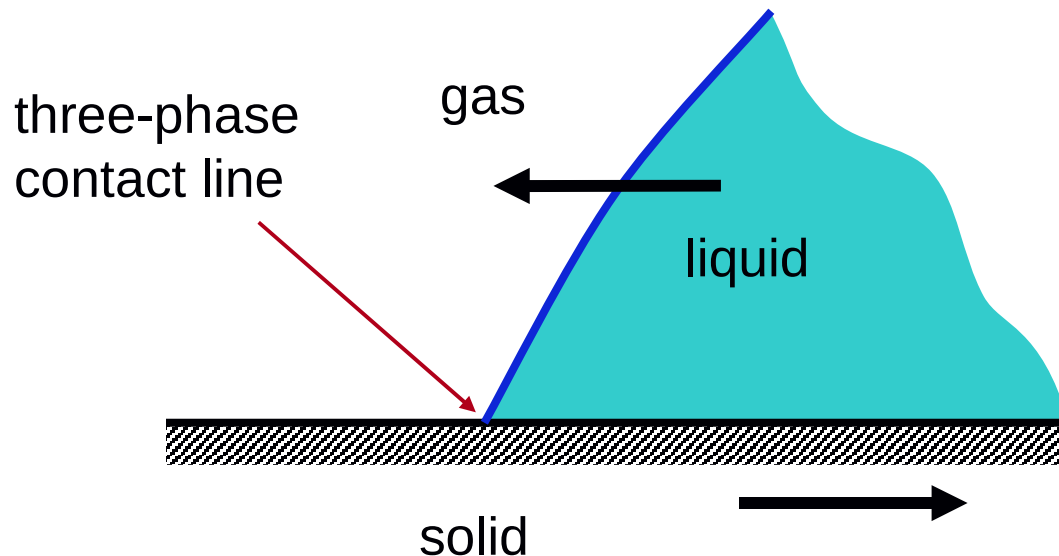
# What are coating flows?



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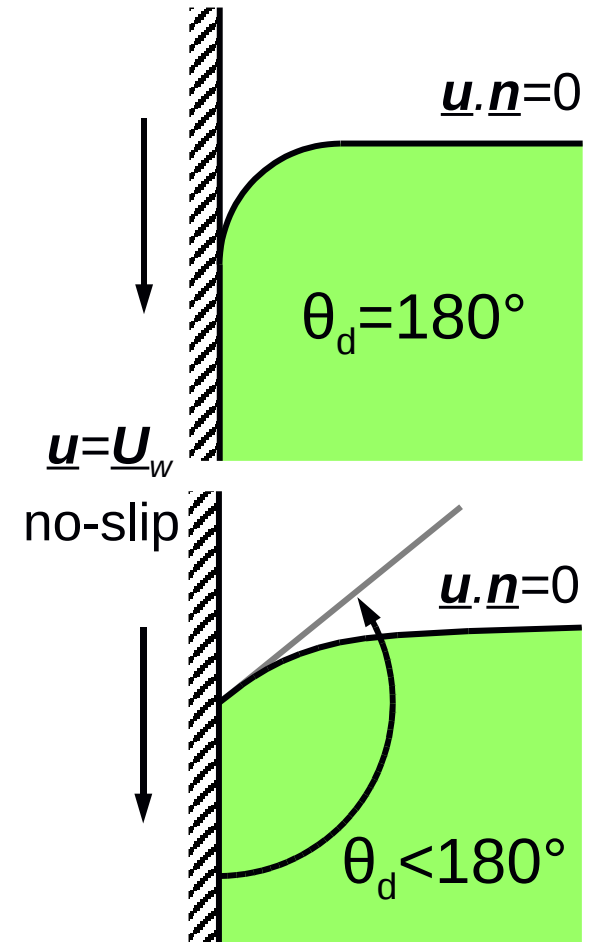


- Can categorize as either **natural wetting** or **forced wetting**
  - Droplet spreading → natural wetting
  - Coating flows → forced wetting
- Central to wetting is the problem of the **moving contact-line**



In first part of talk  
gas will be neglected  
– modelling only  
includes liquid

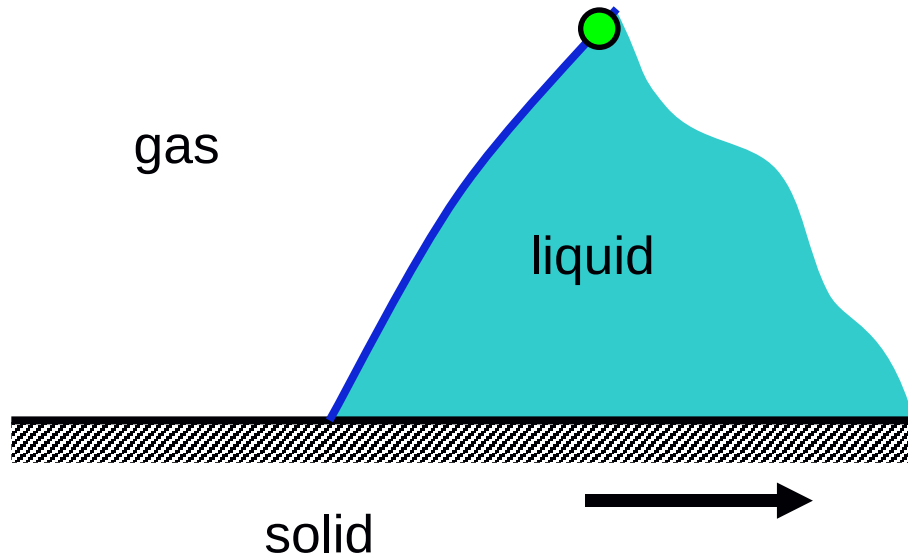
- Presence of a finite contact angle causes problems
- Boundary conditions on liquid-gas and solid-liquid interfaces are in conflict at the contact line
- Contact line is stationary, but boundary is moving
- **Shear stress is infinite**
- Usually have to **prescribe contact angle and slip (relieve stress)**



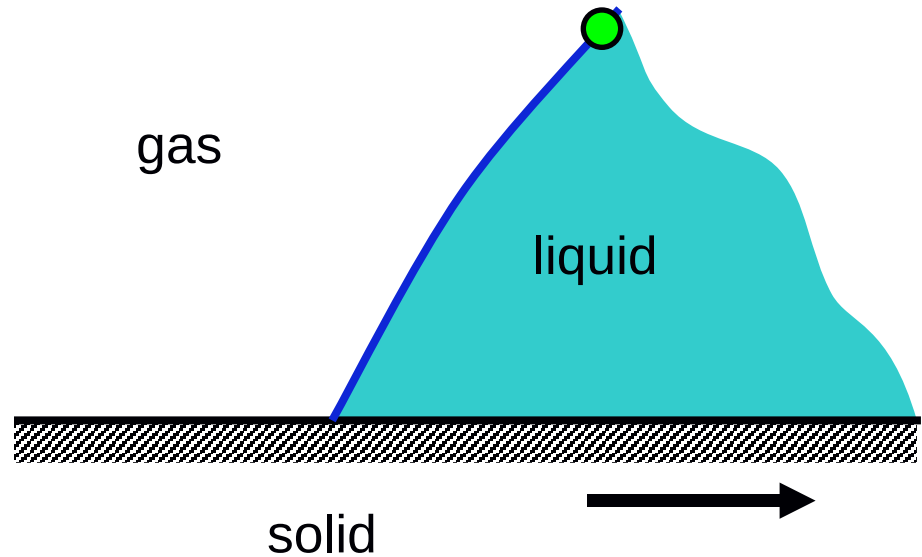
# Problems With Slip Models



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- 'Rolling' motion observed in experiments
- Particle on liquid-gas interface passes through contact line and onto solid-liquid interface

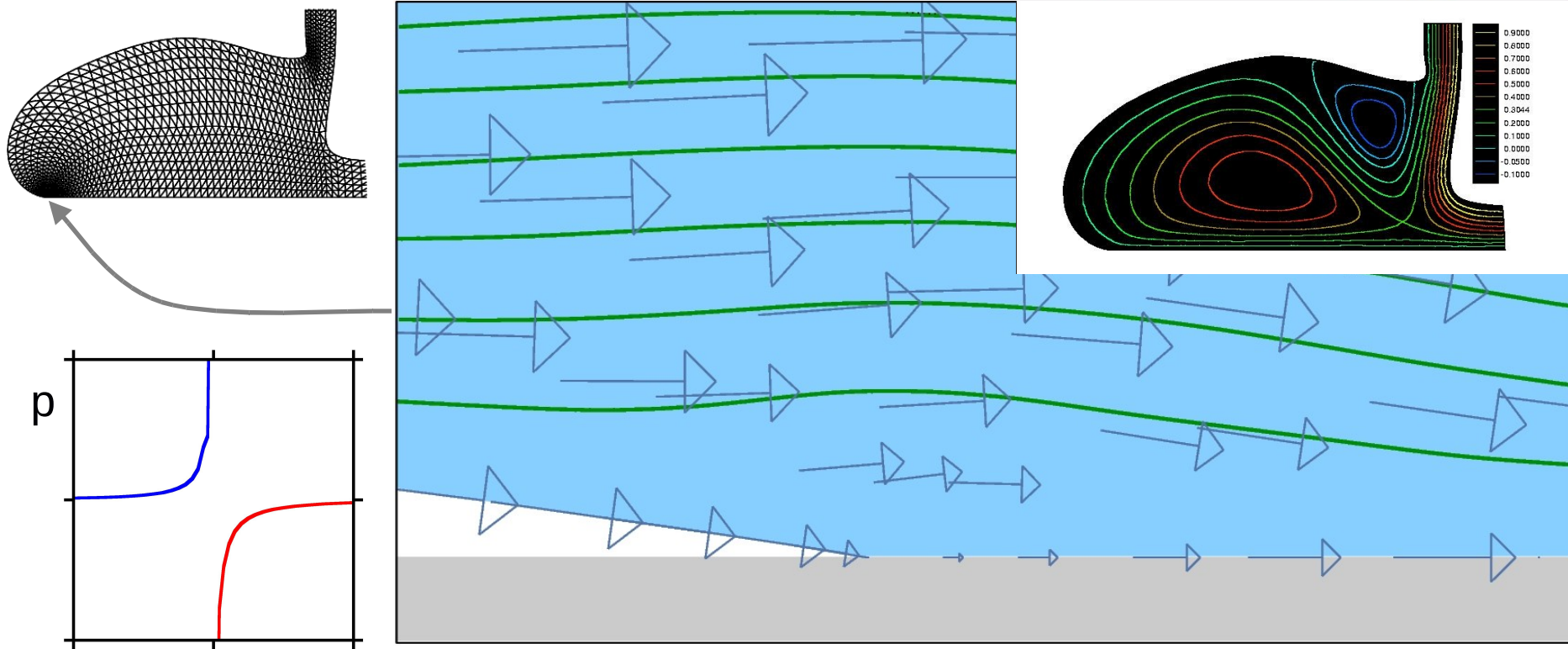


- 'Sliding' motion produced by most slip models
- Particle on liquid-gas interface never reaches contact line, because  $u \rightarrow 0$

# Problems With Slip Models



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- Coating flows driven at high speed often have large  $\theta_d$
- Slip region produces **obstacle-type flow**
- Pressure is **singular**

# `Interface Formation Model' of Shikhmurzaev



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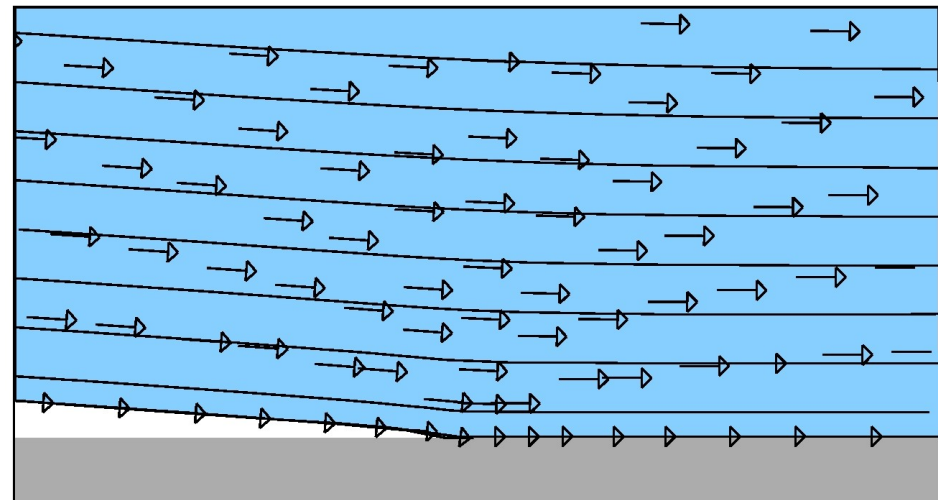
- Based on idea of relaxation of surface tension
  - *Interfacial tension changes smoothly from L-G to L-S*
  - *Near contact line interfacial tensions deviate from equilibrium values*
  - *Force balance at contact line gives contact angle as function of flow*
- The **usual** kinematic and impermeability conditions are replaced with equations describing fluxes between the bulk and the interfaces
- The stress conditions are modified to account for **variable interfacial tensions**
- Liquid velocity at the contact line is not zero – it is determined as part of the solution. Rolling motion preserved.
- Dynamic contact angle is obtained from solution

# `Interface Formation Model' of Shikhmurzaev



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- Model has been used successfully for Stokes flows
  - Lukyanov & Shikhmurzaev (2007) *Phys. Rev. E* **75**, 051604
  - considered a microfluidic curtain coater
  - observed variation of  $\theta_d$  with a number of flow parameters
  - used a combined finite element-boundary integral element method
- Navier-Stokes **finite element** solutions for the full-scale curtain coater are now possible.
- **But**, air-entrainment predictions are not possible.







- Supported by molecular dynamics simulations, a diffuse interface for the **liquid-gas**, **solid-liquid** and **solid-gas** is more amenable to varying interfacial density.
- Diffuse interfaces can rupture and so could help to predict the important aspect of wetting failure, i.e. air-entrainment.
- Several multiphase lattice Boltzmann (LB) approaches exist.
- Wetting line tests for an LB method are:
  - 1. Forced wetting with failure
  - 2. Wetting line hysteresis
  - 3. Natural wetting (spreading/sticking) [agreement with experiments]

Based on work of He, Chen & Zhang (1999)

Use mean-field approximation for intermolecular attractions, and include an exclusion volume effect to...

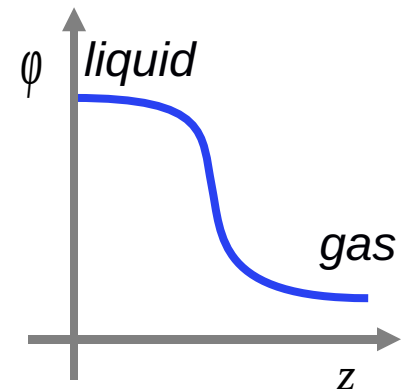
Rework force term in Boltzmann equation into a surface tension force

Use **non-ideal equation of state** to achieve phase separation

Introduce an **index function**,  $\phi$ , to track the interface between two phases

Results in a **diffuse interface** model

- *index function, and fluid density, changes smoothly but rapidly between phases*



# Finite-density Multiphase Lattice Boltzmann Equations



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- Following **He, Chen and Zhang**, two LB equations with forcing are derived for  $f_i$  and  $g_i$ , the moments of which give the macroscopic properties; mass and momentum densities, and pressure respectively as

$$\phi(x_\alpha, t) = \sum f_i(x_\alpha, t) \quad (\phi \text{ tracks density} = \text{index function})$$

$$\rho u_\alpha(x_\alpha, t) = \frac{1}{RT} \sum e_{i\alpha} g_i(x_\alpha, t) + \frac{1}{2} \left[ \kappa \rho \frac{\partial}{\partial x_\alpha} (\nabla^2 \rho) + B_\alpha \right] \delta t$$

$$p(x_\alpha, t) = \sum g_i(x_\alpha, t) - \frac{1}{2} u_\alpha \frac{\partial \psi(\rho)}{\partial x_\alpha} \delta t$$

# Finite-density Multiphase Lattice Boltzmann Equations



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- Model is for a liquid and its gas. Values of the index function for liquid,  $\phi_L$ , and gas,  $\phi_G$  are obtained from the EoS and Maxwell's equal area construct.
- The liquid and gas values of  $\phi$  can be used to account for the different fluid properties between the phases, that is

$$\rho(\phi) = \rho_G + \frac{\phi - \phi_G}{\phi_L - \phi_G} (\rho_L - \rho_G)$$

- The same is true for the **viscosity,  $\mu$** .
- Can be applied with MRT (see Premnath and Abraham (2007))

Use the approach of Iwahara et al. (2003)

Define a **surface affinity** – a normalised surface density

$$\alpha_s = \frac{\phi - \bar{\phi}}{\phi_L - \bar{\phi}}, \quad \text{where } \{ \bar{\phi} = (\phi_L + \phi_G)/2 \}$$

A **planar** interface has the profile (Rowlinson & Widom 1982)

$$\phi(z) = \bar{\phi} - \frac{1}{2}(\phi_L - \phi_G) \tanh\left(\frac{z - z_0}{\delta}\right)$$

The liquid-gas surface tension is therefore

$$\sigma_{LG} = K \int_{-\infty}^{\infty} \left( \frac{\partial \phi}{\partial z} \right)^2 dz = \frac{K(\phi_L - \phi_G)^2}{4\delta} \int_{-1}^1 (1 - \alpha^2) d\alpha$$

$$\frac{K(\phi_L - \phi_G)^2}{3\delta}$$

Similar expressions for the solid-liquid and solid-vapour surface tensions substituted into Young's equation give

$$\cos \theta_s = \alpha_s (3 - \alpha_s^2) / 2$$

Static contact angle can be specified via the surface affinity

Index-function density at boundary given by

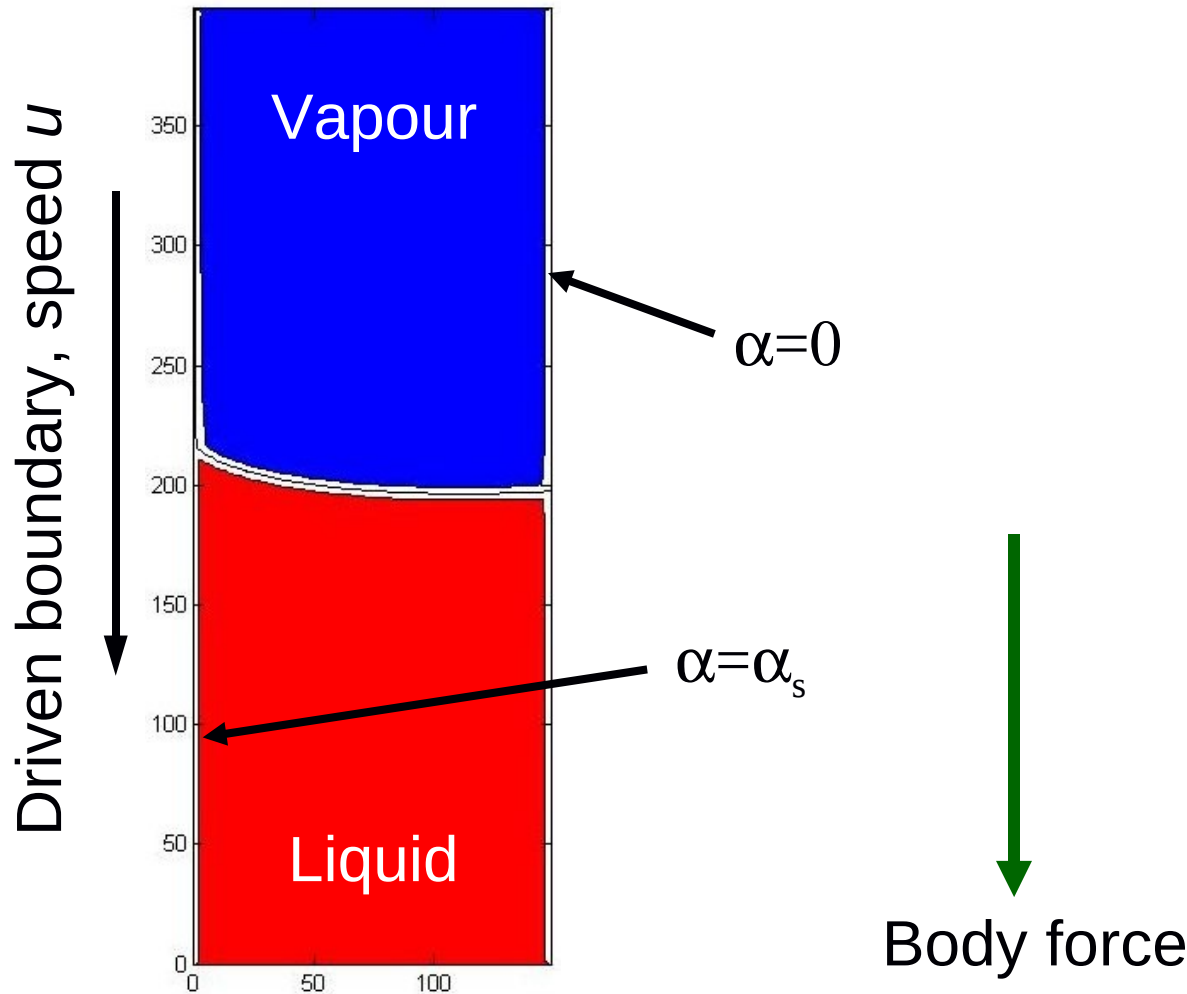
$$\varphi = \bar{\phi} + (\phi_L - \bar{\phi}) \alpha_s, \quad \{ \bar{\phi} = \frac{1}{2} (\phi_L + \phi_G), \quad -1 \leq \alpha_s \leq 1 \}$$

# Test Problem 1 (Forced wetting to failure)



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- Two-phase cavity
- Solid walls

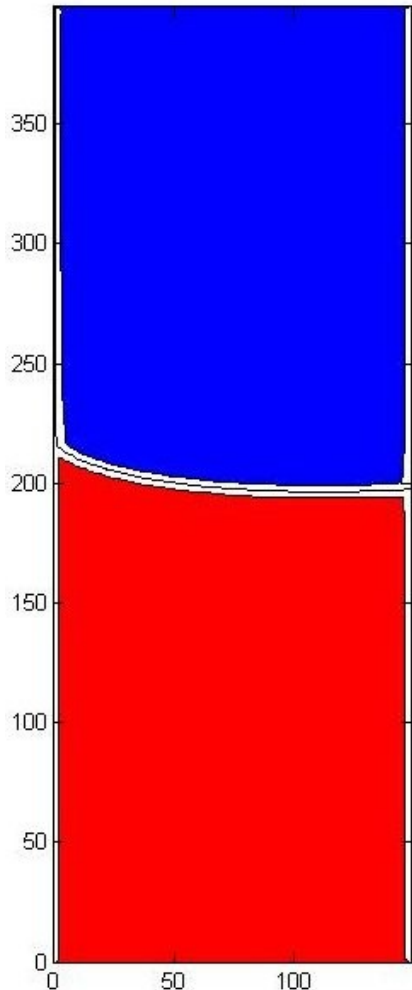


# Test Problem 1 (Forced wetting to failure)



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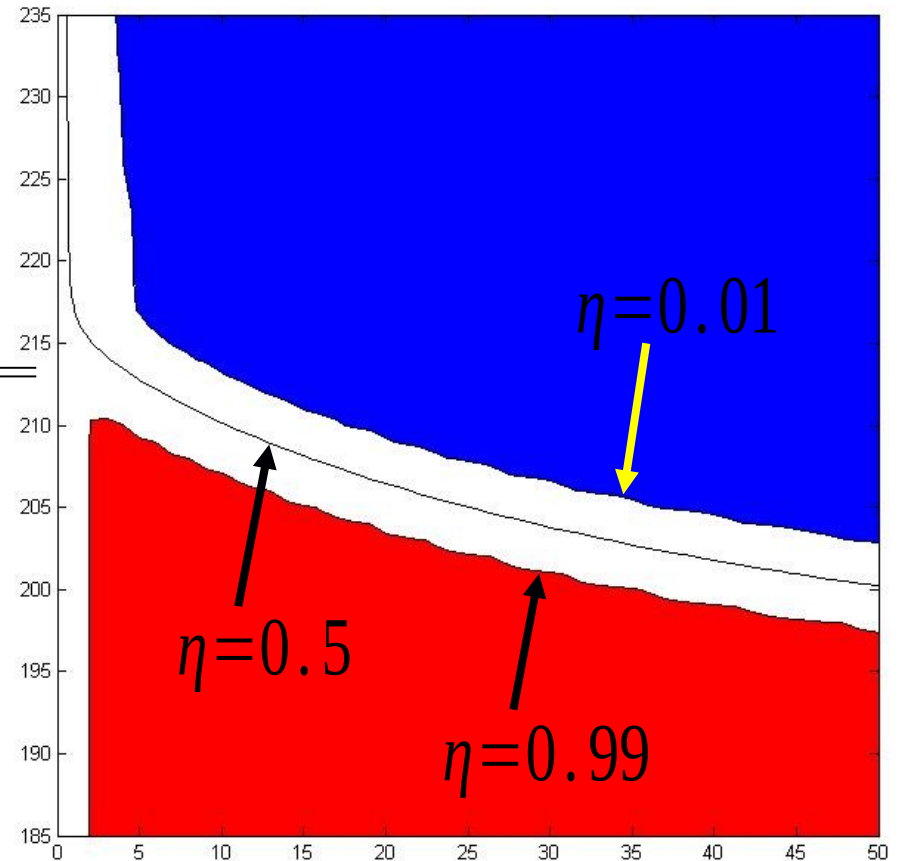
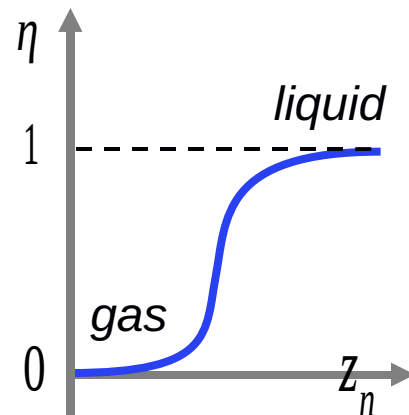
## Static Case – Interface Shape



$$\alpha_s = 0.4$$

$$\Rightarrow \theta_s = 55^\circ$$

$$\eta = \frac{\phi - \phi_G}{\phi_L - \phi_G} \eta =$$



**Where is the contact 'line'?**  
**Which is 'the' contact angle?**

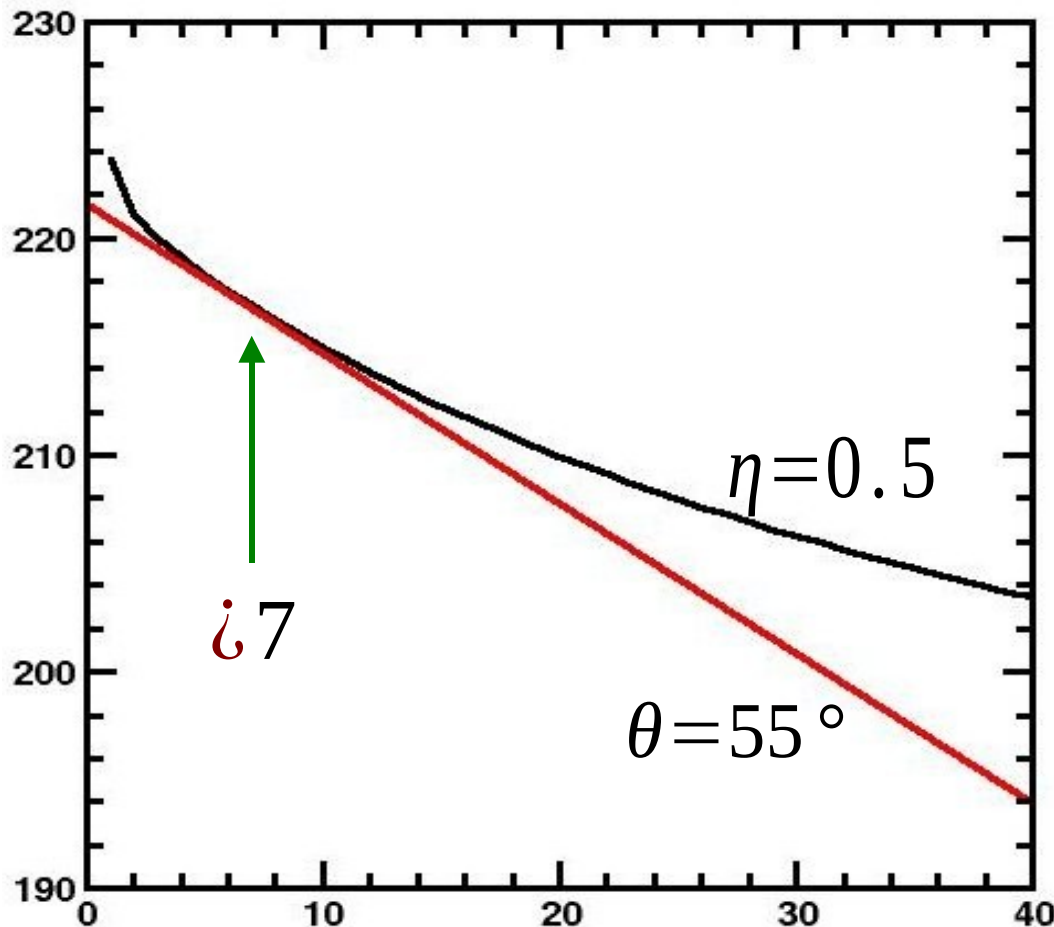


# Test Problem 1 (Forced wetting to failure)



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## Static Case – Contact Angle



Angle of interface matches imposed angle at roughly 7 lattice units away from the boundary

Just outside the diffuse three-phase contact region

Use this as the point to measure variation in contact angle

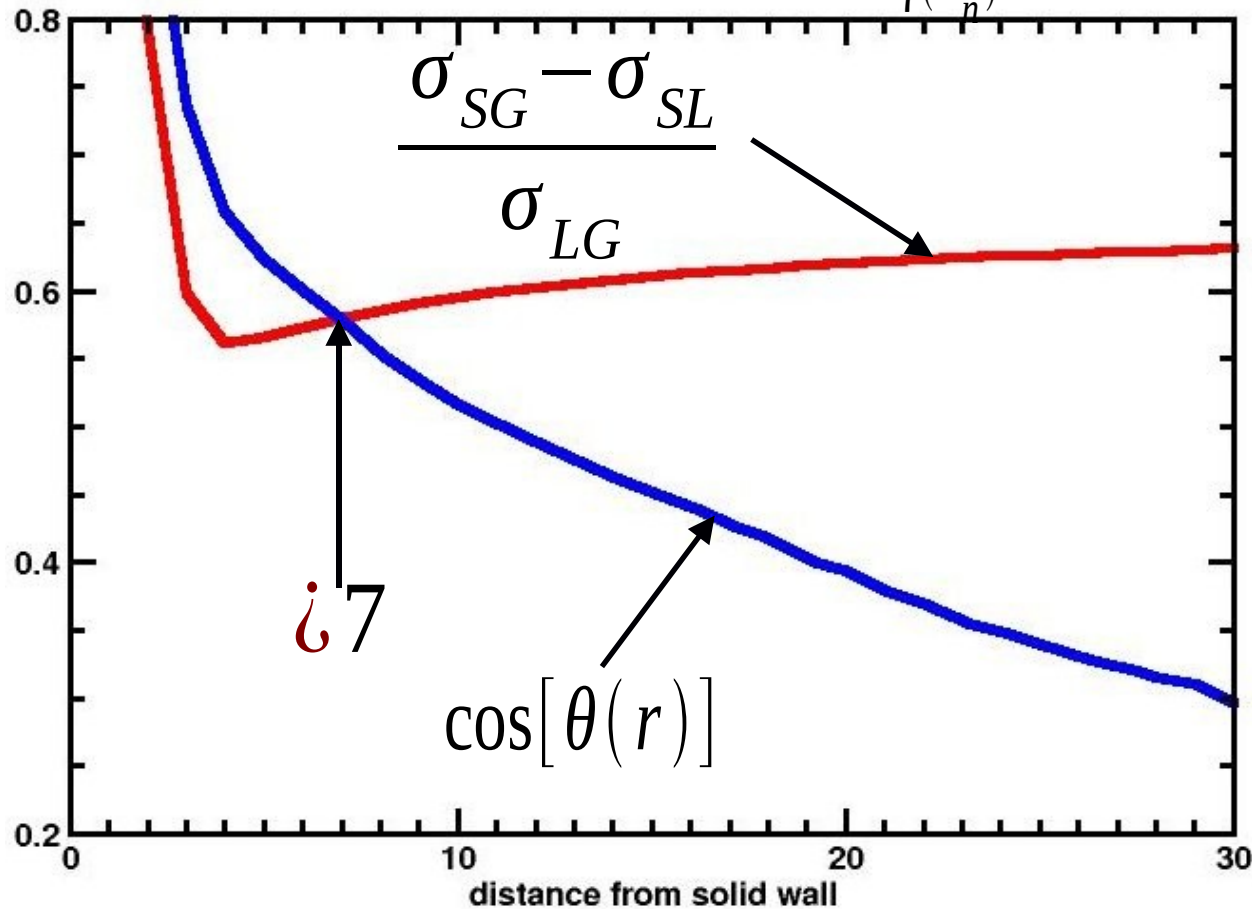
# Test Problem 1 (Forced wetting to failure)



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Static Case – Young equation

$$\sigma = K \int_{\eta(z_n)=0.01}^{\eta(z_n)=0.99} [\nabla \phi]^2 dz_n$$



# Test Problem 1

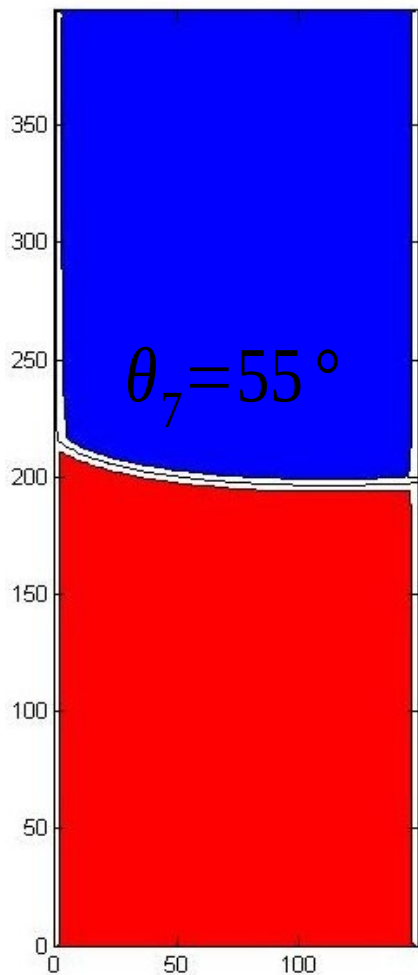
## (Forced wetting to failure)



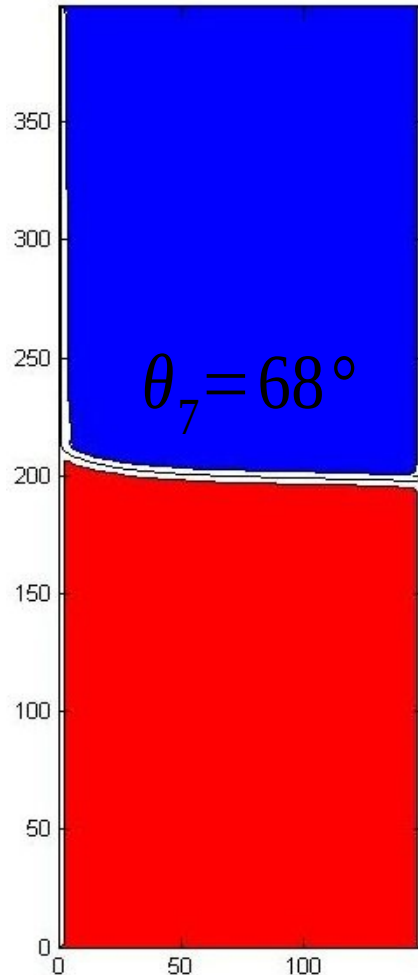
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Dynamic Case – Contact Angle

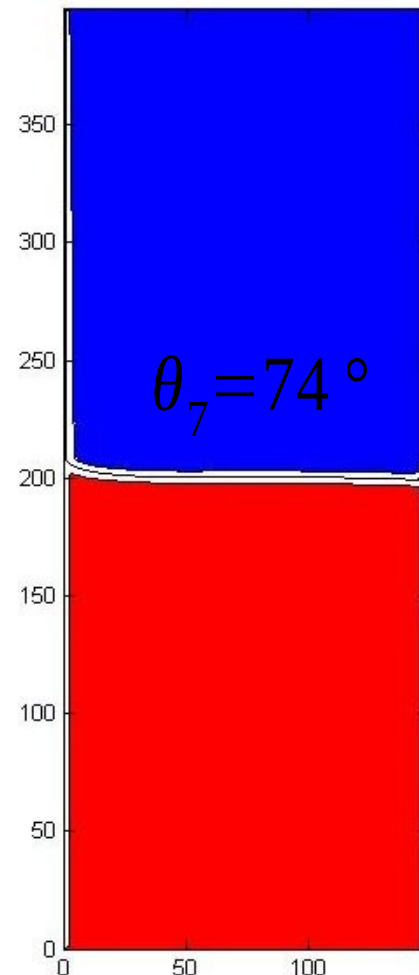
$$u=0$$



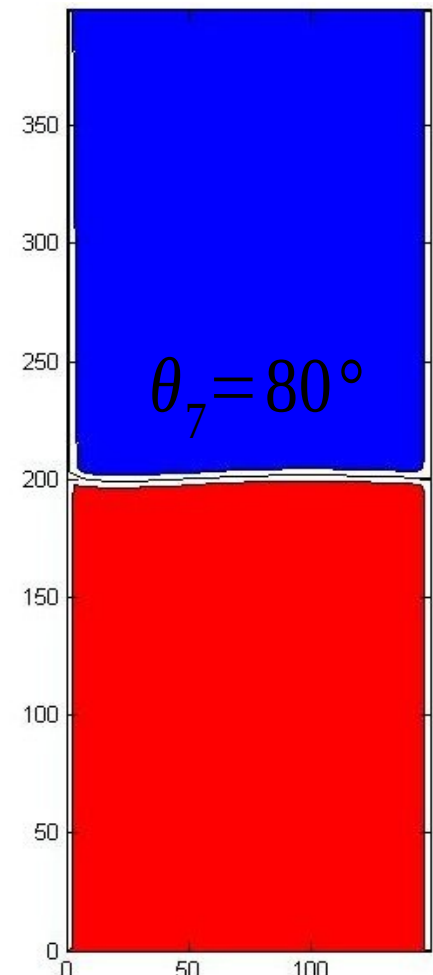
$$u=1 \times 10^{-3}$$



$$u=1.5 \times 10^{-3}$$



$$u=2 \times 10^{-3}$$

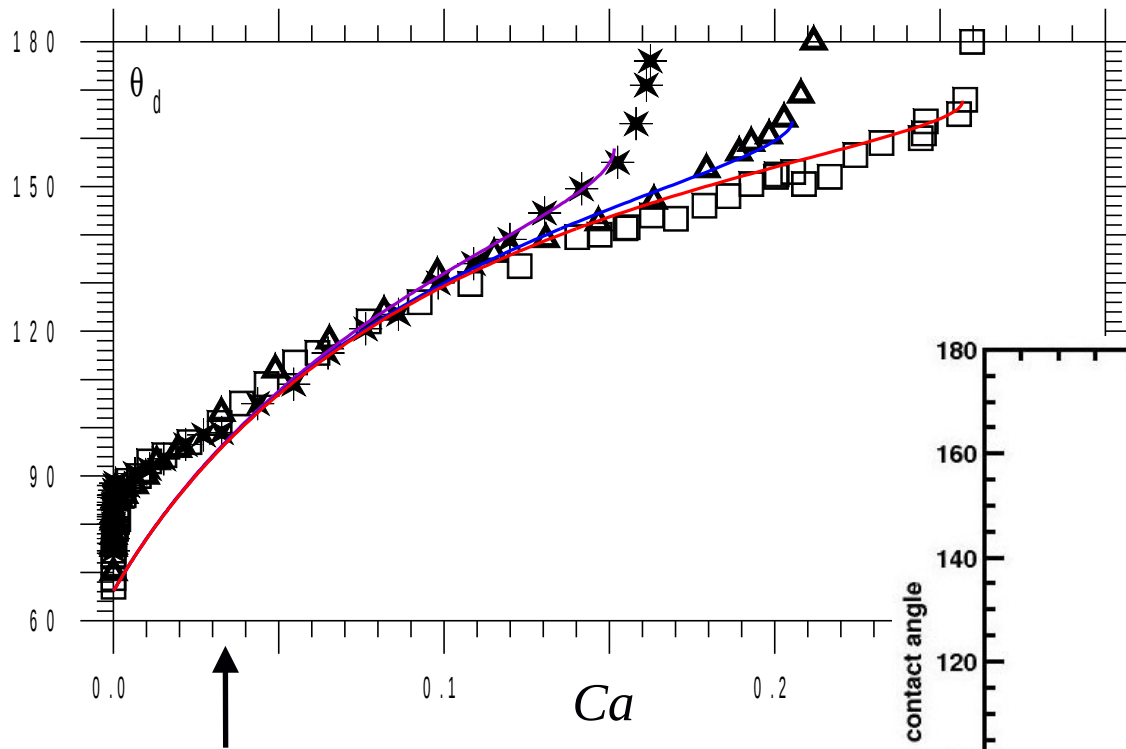


# Test Problem 1 (Forced wetting to failure)



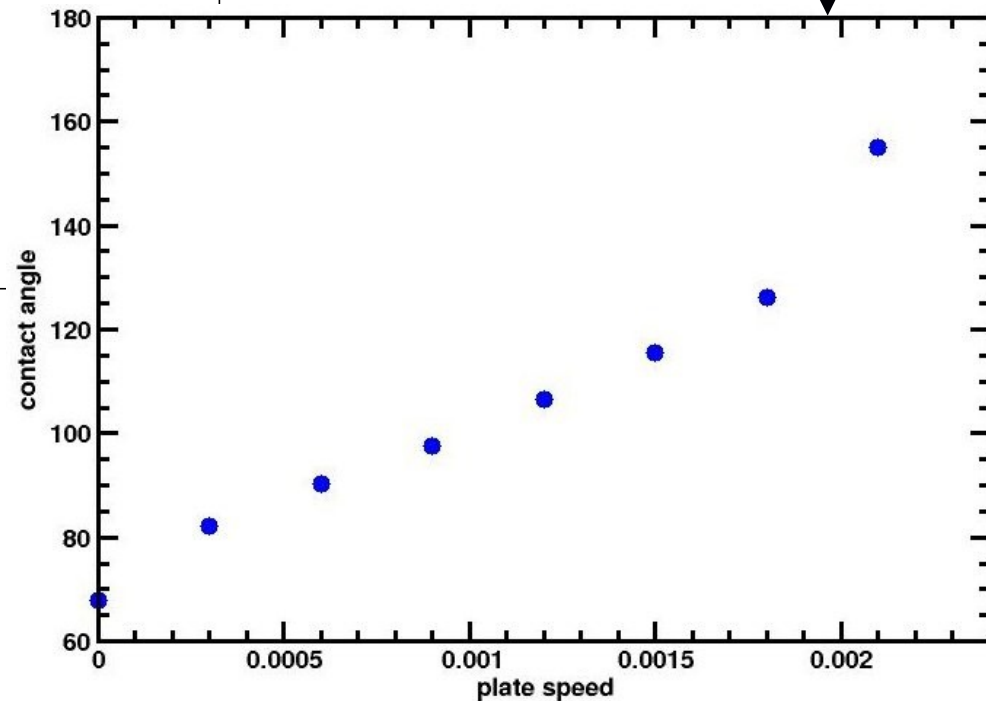
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## Contact Angle Versus Speed



Experiments of Blake *et al.*  
*J. Colloid Interf. Sci.* **253**,196 (2002)

Lattice Boltzmann results

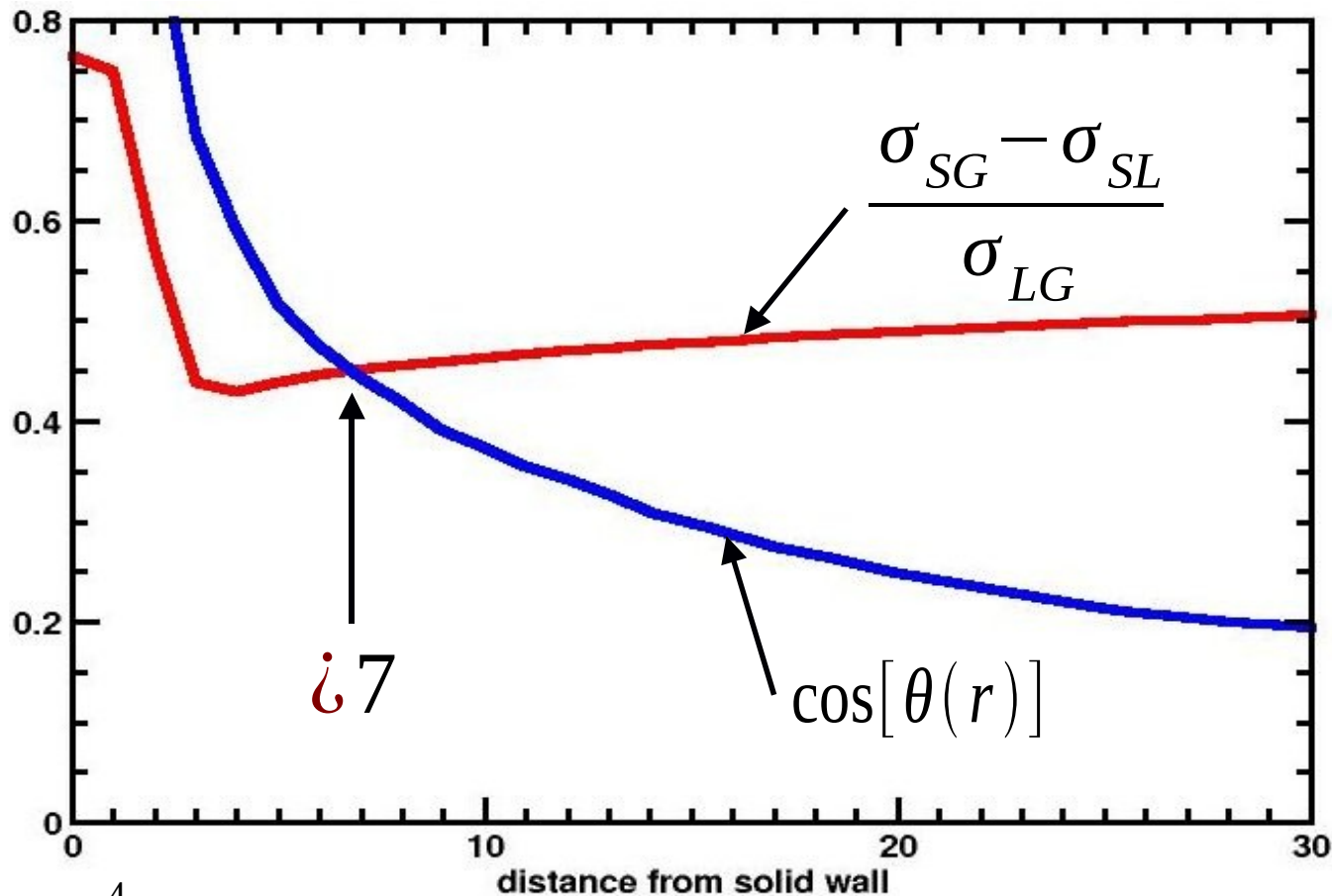


# Test Problem 1 (Forced wetting to failure)



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Dynamic Case – Young Equation



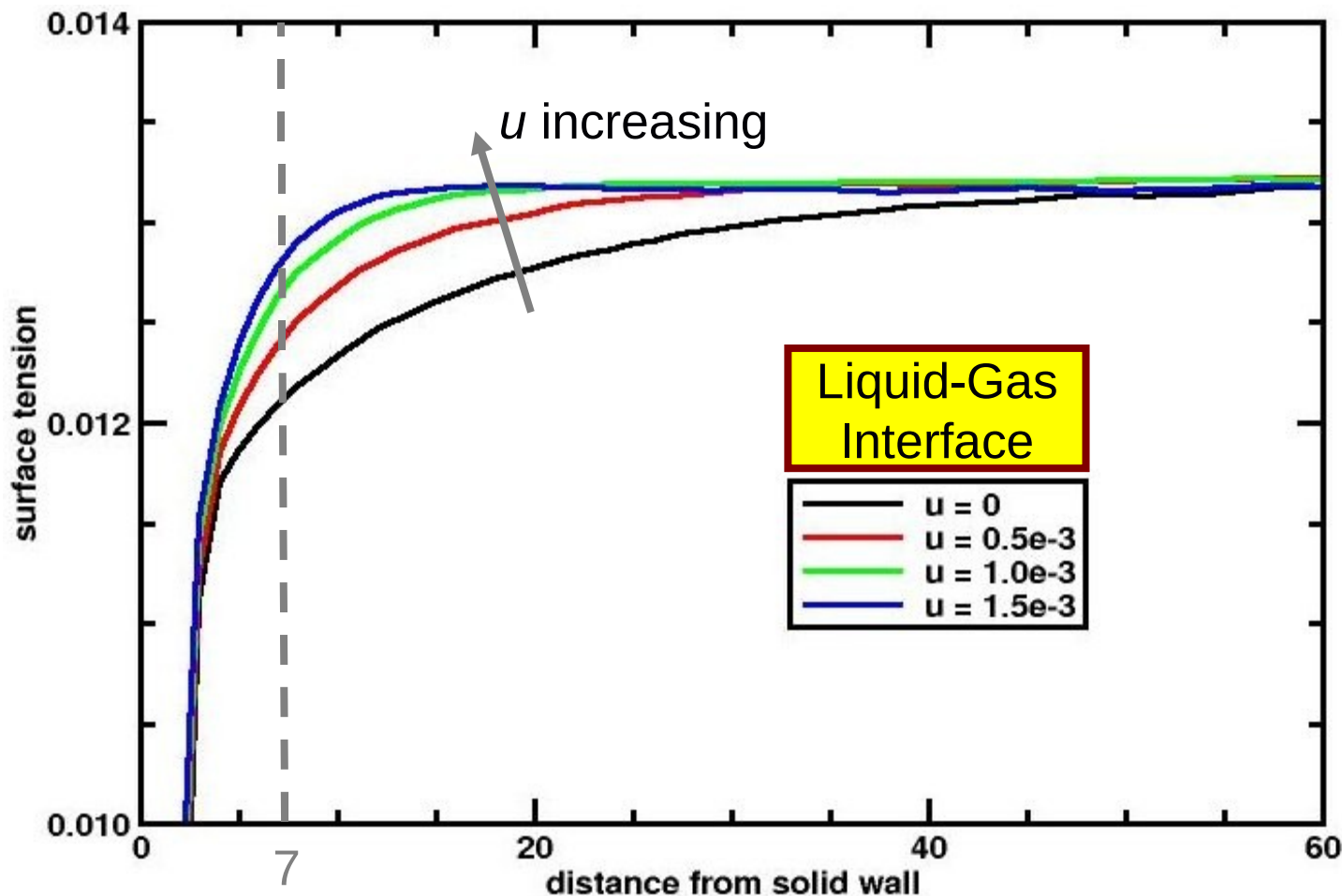
$$u = 5 \times 10^{-4}$$

# Test Problem 1 (Forced wetting to failure)



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## Dynamic Case – Surface Tension

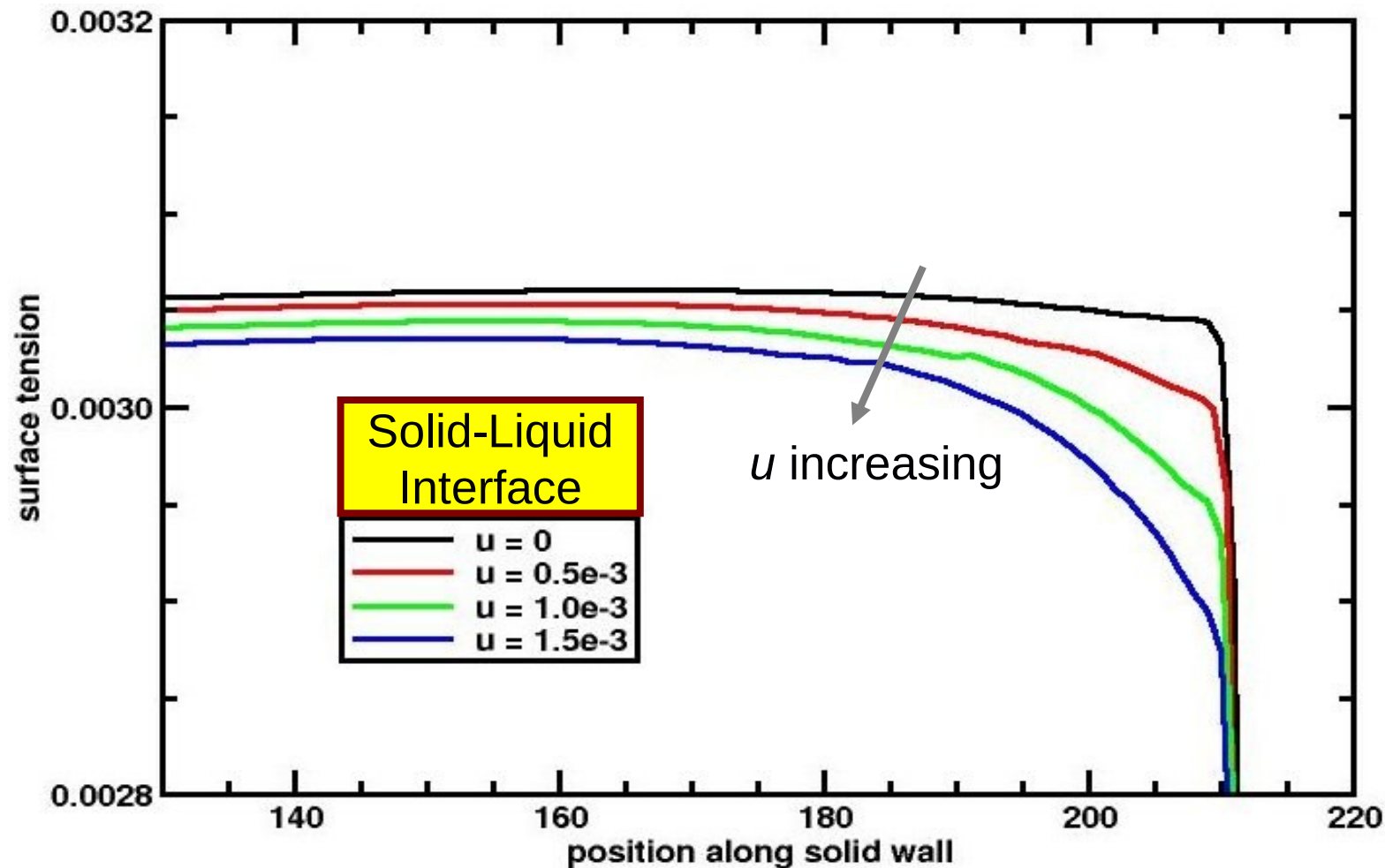


# Test Problem 1 (Forced wetting to failure)



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## Dynamic Case – Surface Energy

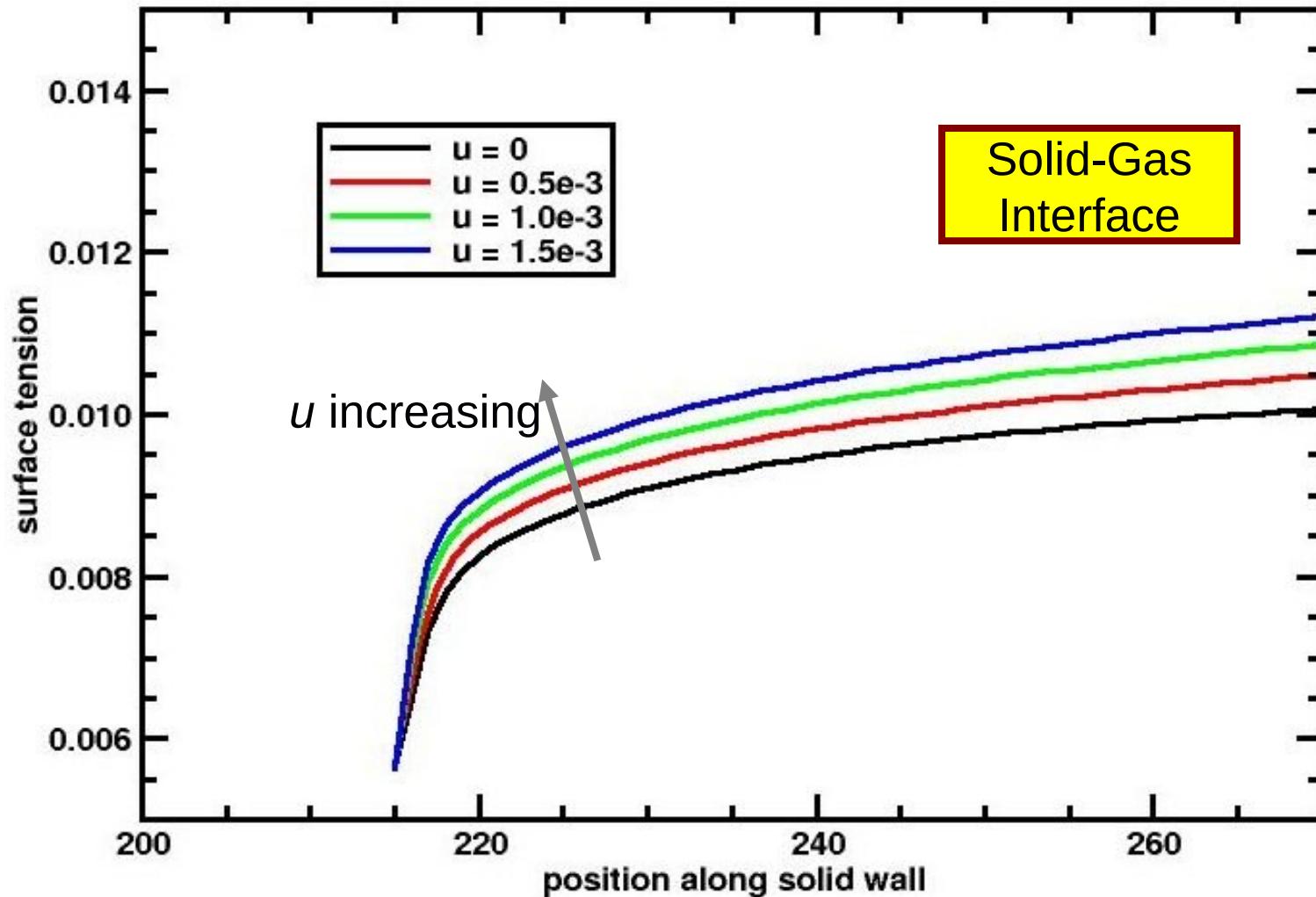


# Test Problem 1 (Forced wetting to failure)



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## Dynamic Case – Surface Energy

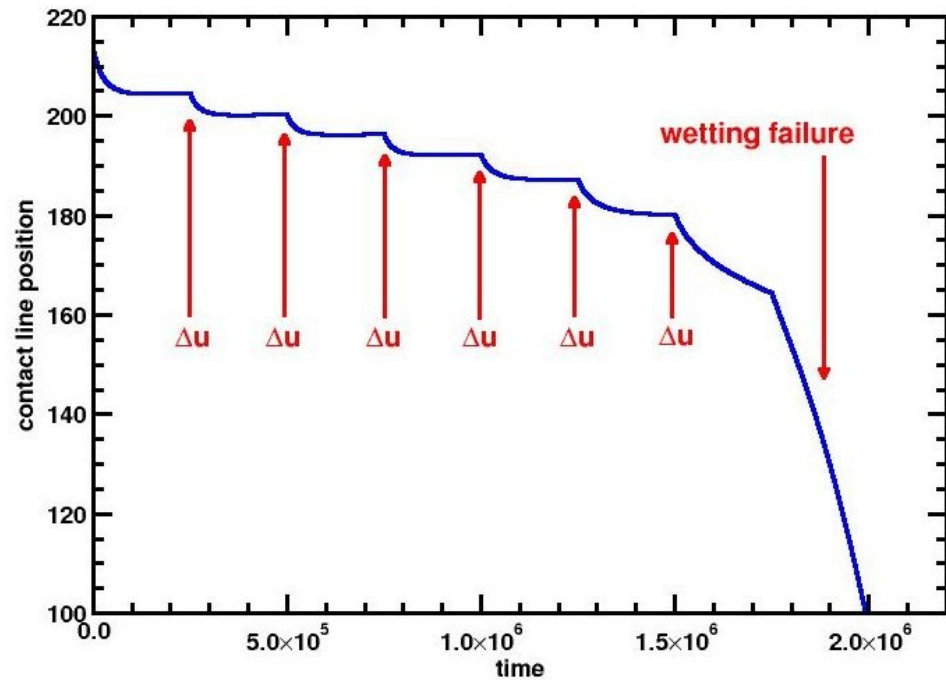
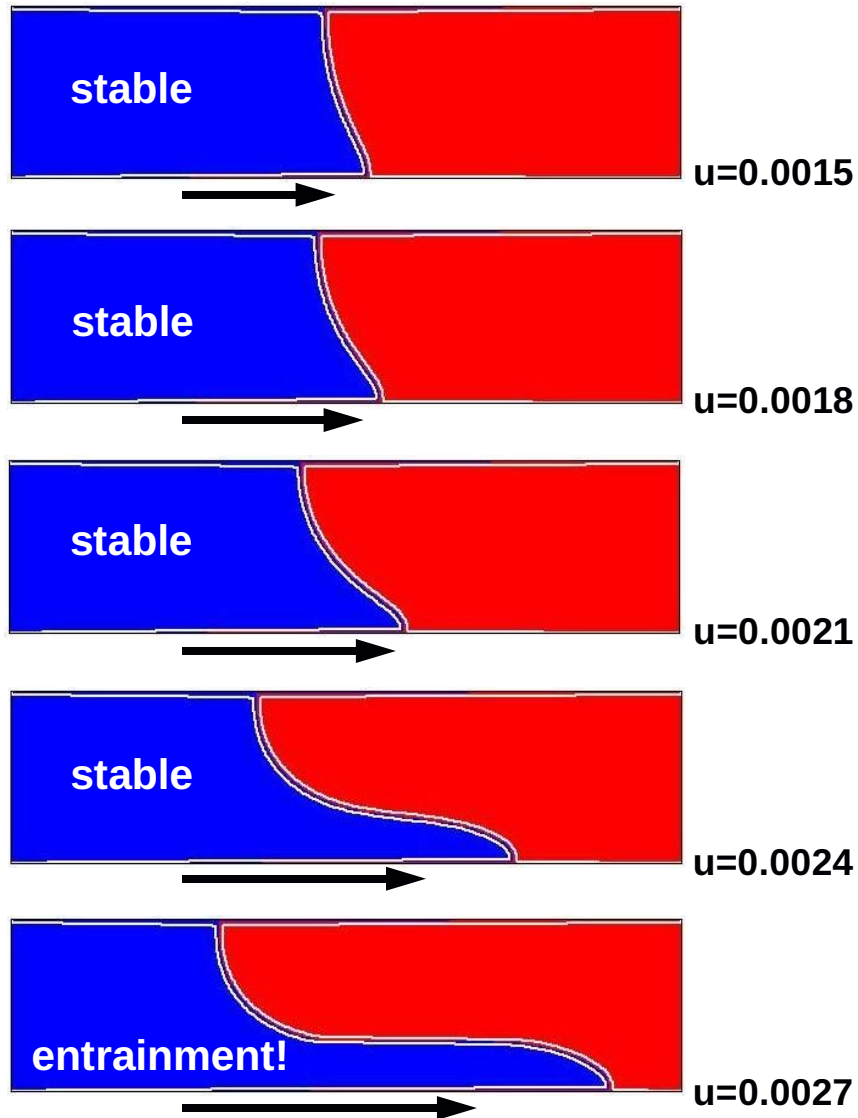




# Test Problem 1 (Forced wetting to failure)



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light phase

dense phase



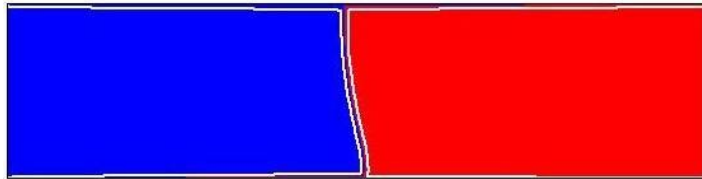
body force (gravity)

# Test Problem 2 (Wetting line hysteresis)

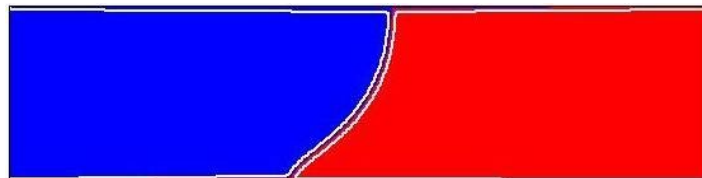
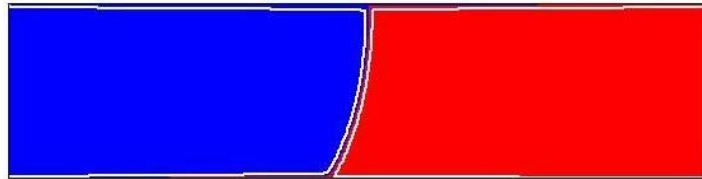
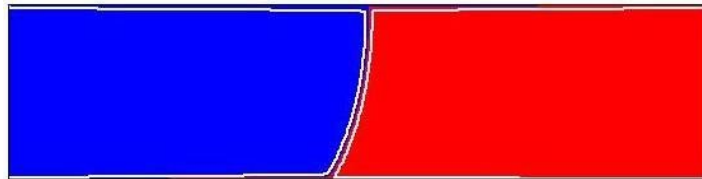


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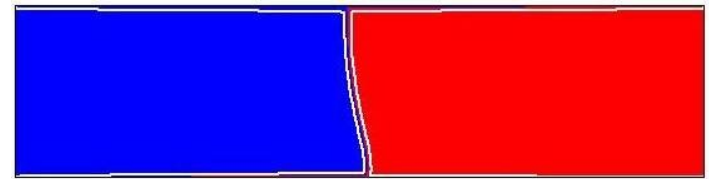
uniform wettability



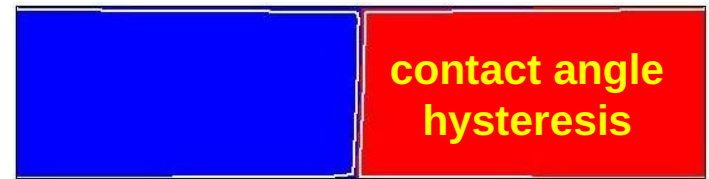
$u > 0$



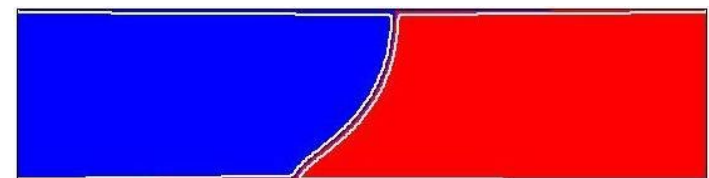
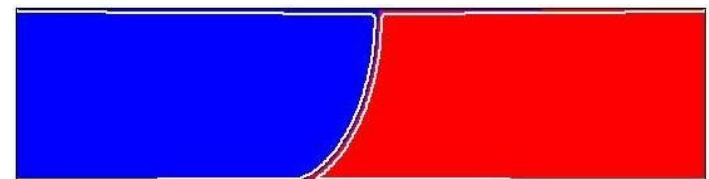
sinusoidally varying wettability



$u > 0$



$u = 0$



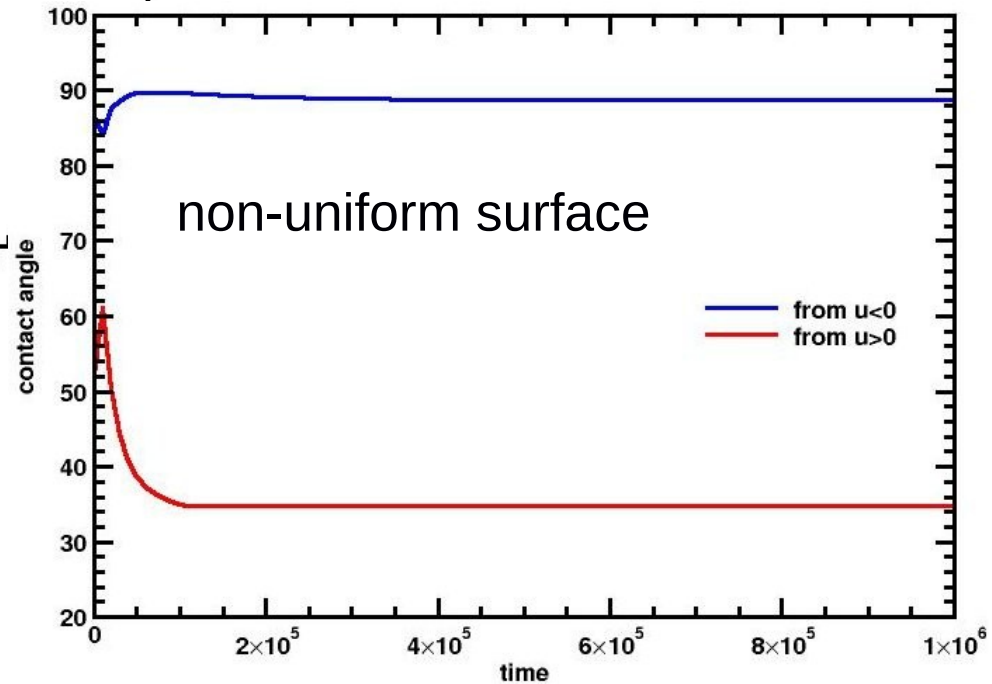
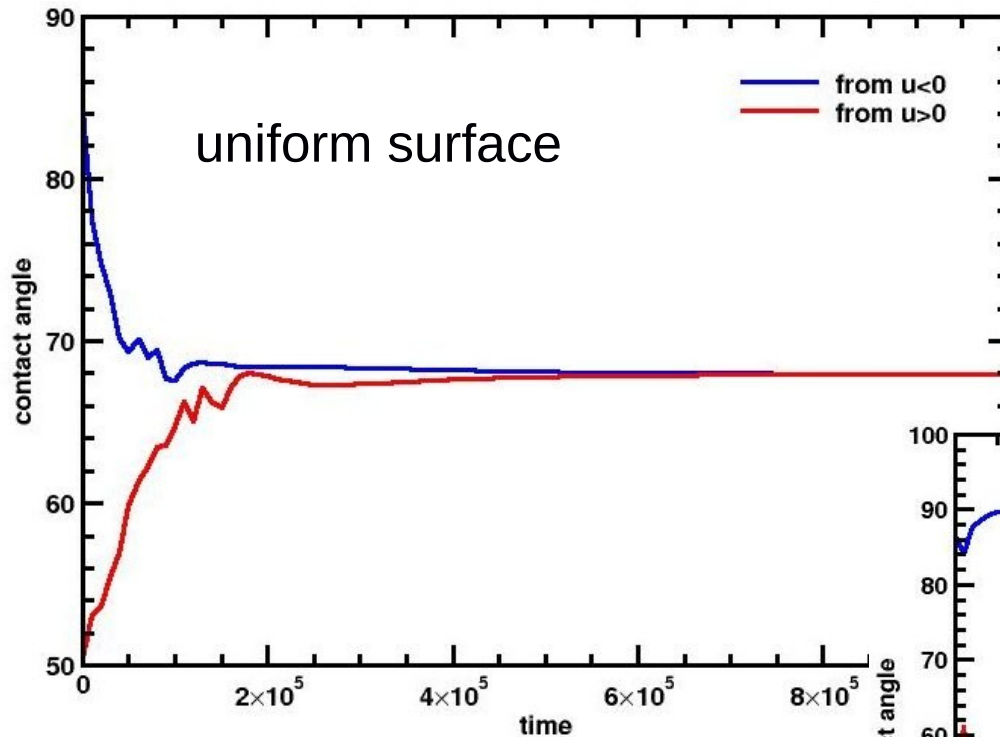
$u < 0$

contact angle hysteresis

# Test Problem 2 (Wetting line hysteresis)



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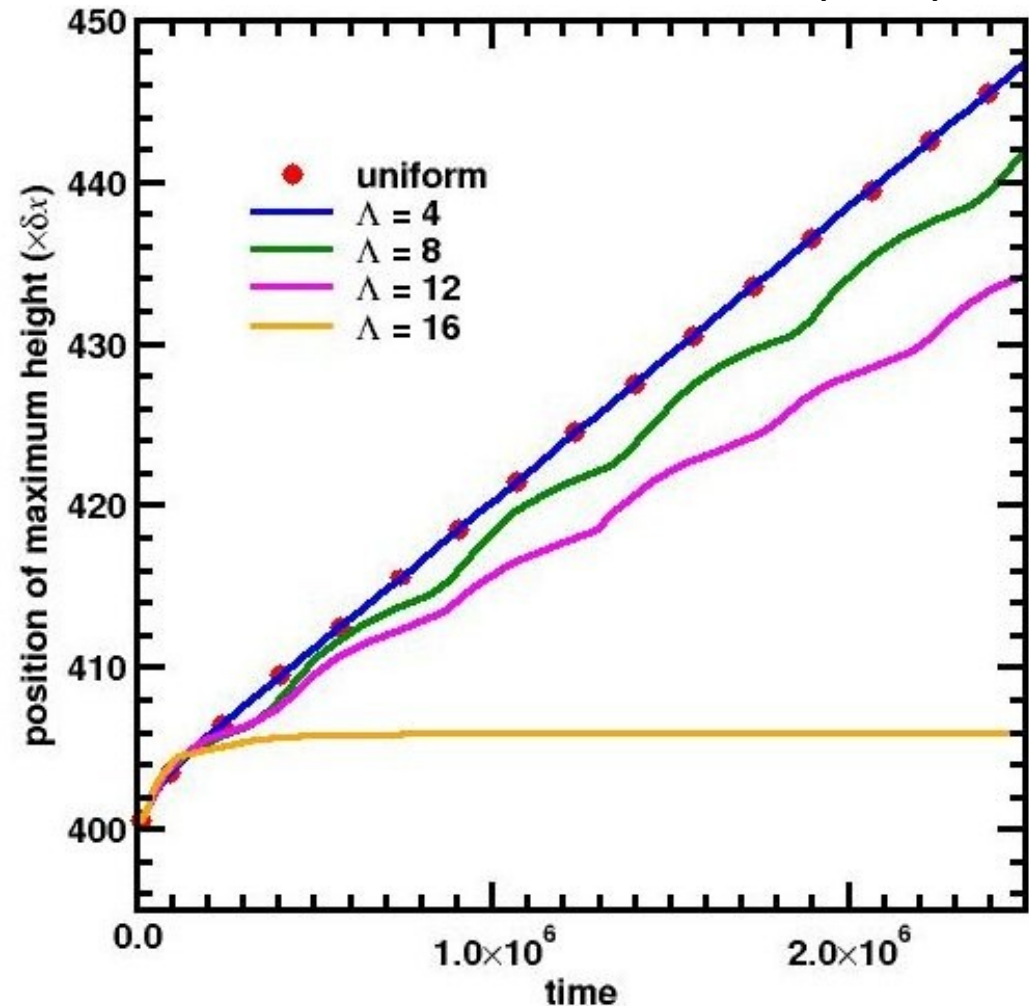
# Test Problem 3 (Natural wetting spreading/sticking)



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- Flow of a droplet down an incline with a sinusoidally varying wettability of wavelength  $\Lambda$
- Varying  $\Lambda$  for fixed interface thickness
- Droplet is pinned for certain values

*Davies, Summers & Wilson (2006)*

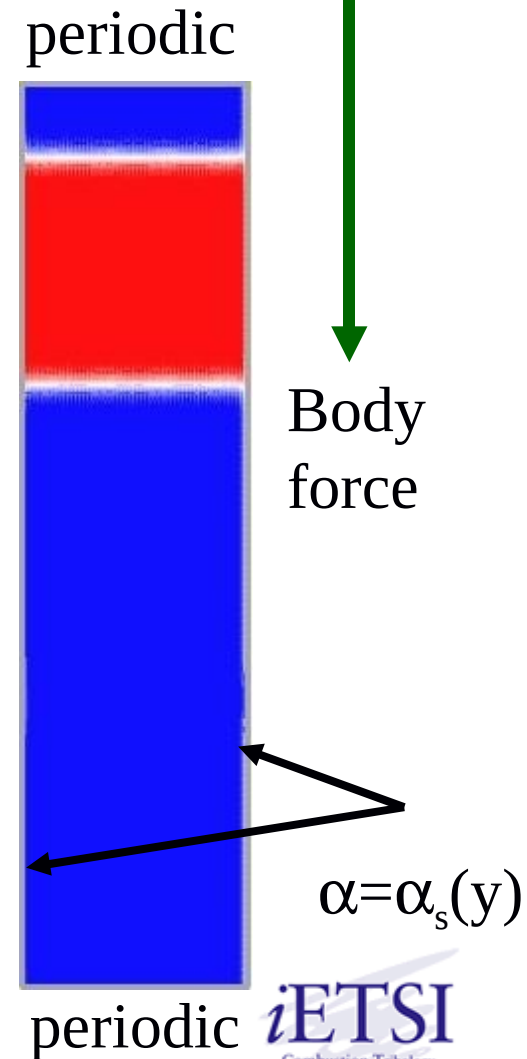


# A Tale of Two Length Scales



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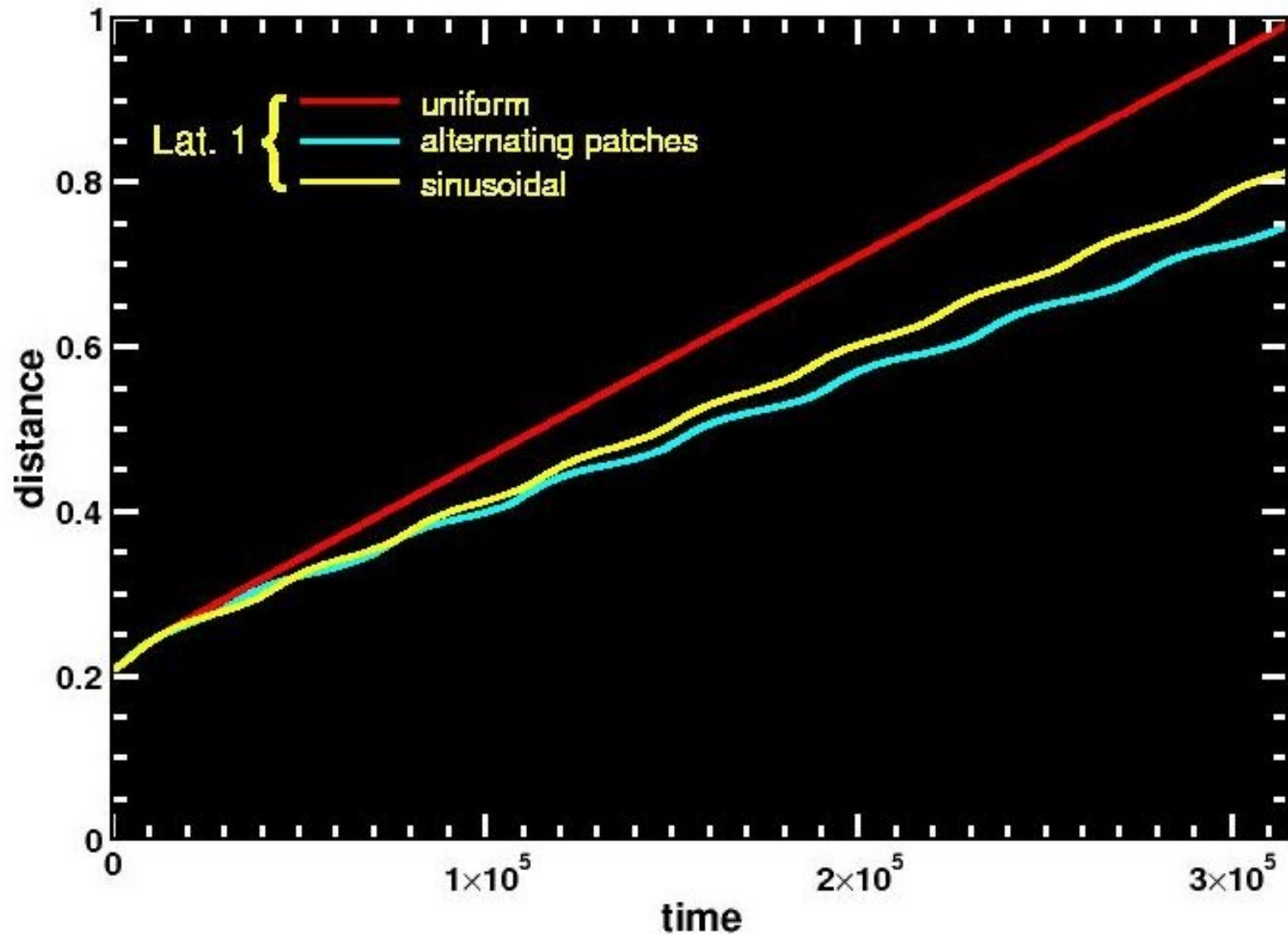
- What is the effect of interface thickness for wetting of non-uniform surfaces?
  - interface thickness versus characteristic size of non-uniformity
- Interface thickness is always  $\sim 4$  or  $5$  lattice units...can scale up problem
- Use two lattices – one twice the size of the other (i.e. twice as dense)
  - need to adjust relaxation time and surface tension parameter to match physical scales on each lattice
- Use sinusoidal/alternating surface affinity



# Motion of slug centre



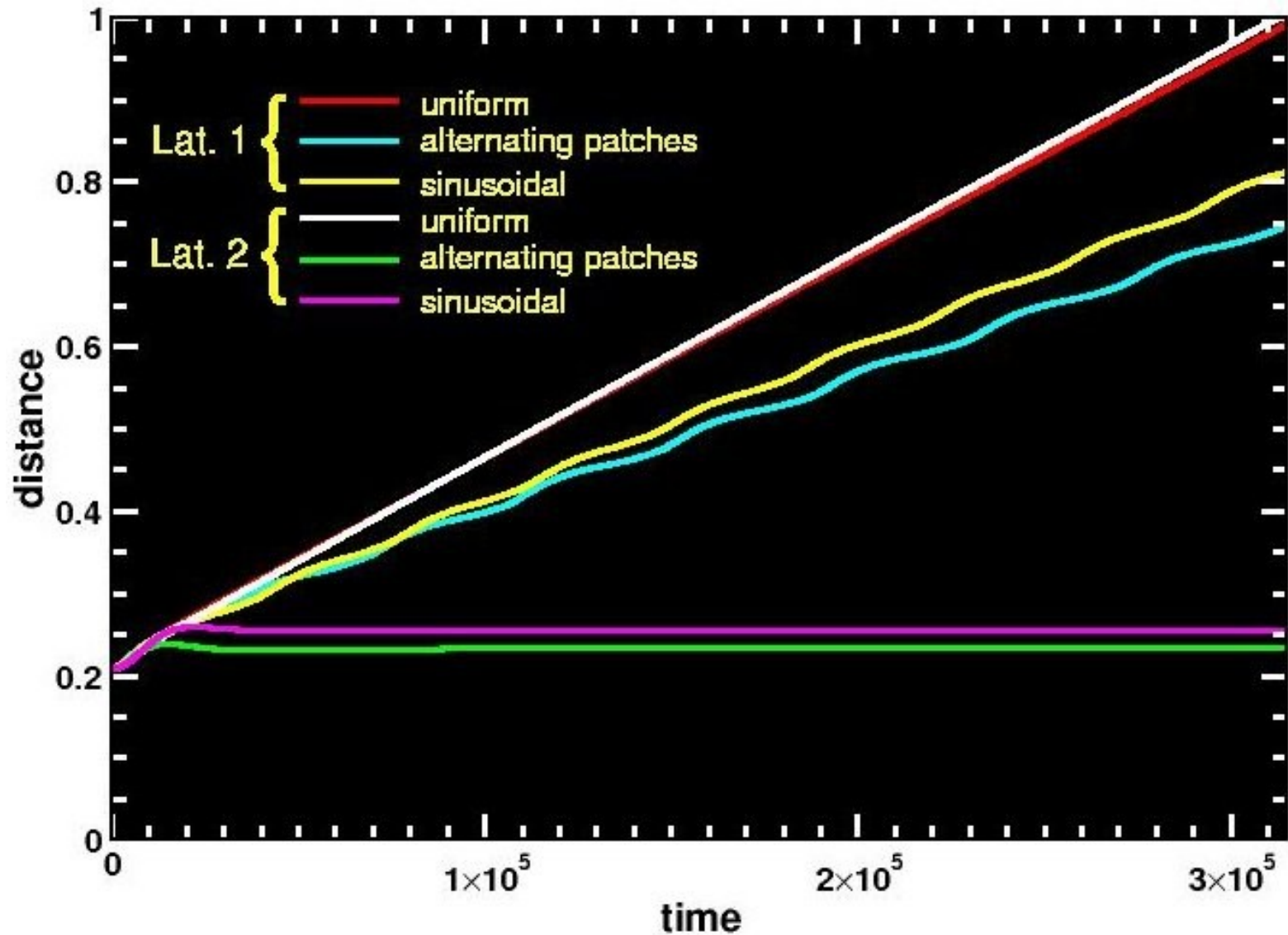
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# Motion of slug centre



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- Density ratio limited
  - *several models now available addressing this issue, though wetting is still an issue for many*
- Issues with the surface tension and surface energies.
  - *Calculations of surface tension/energy via the thermodynamic (Cahn) approach is for static conditions, can we define a “mechanical” approach for the dynamic situation.*
- Wetting models for moving rough boundaries needed
  - *interface thickness a key factor*



- Able to capture qualitatively many wetting phenomena (static contact angle, forced wetting to failure, contact line hysteresis and natural wetting to the point of sticking on a non-uniform surface)
- Simple model and algorithm – only one wetting parameter
- Care needed in understanding effect of interface thickness (as this will dictate the length scale?)
- Work needed to make quantitatively accurate (is the thermodynamic description of surface energy sufficient?)